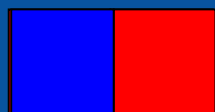
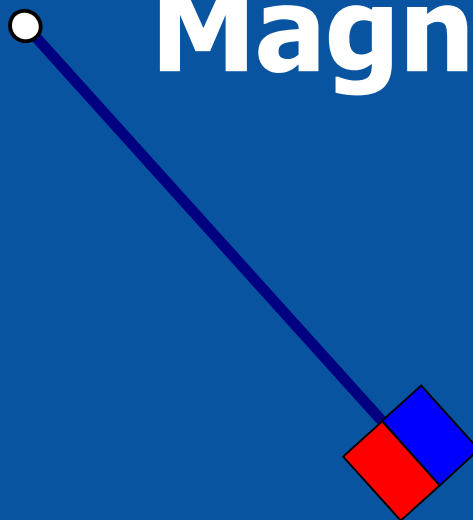




Russia IYPT

Magnetic pendulum

Vitaly Matiunin

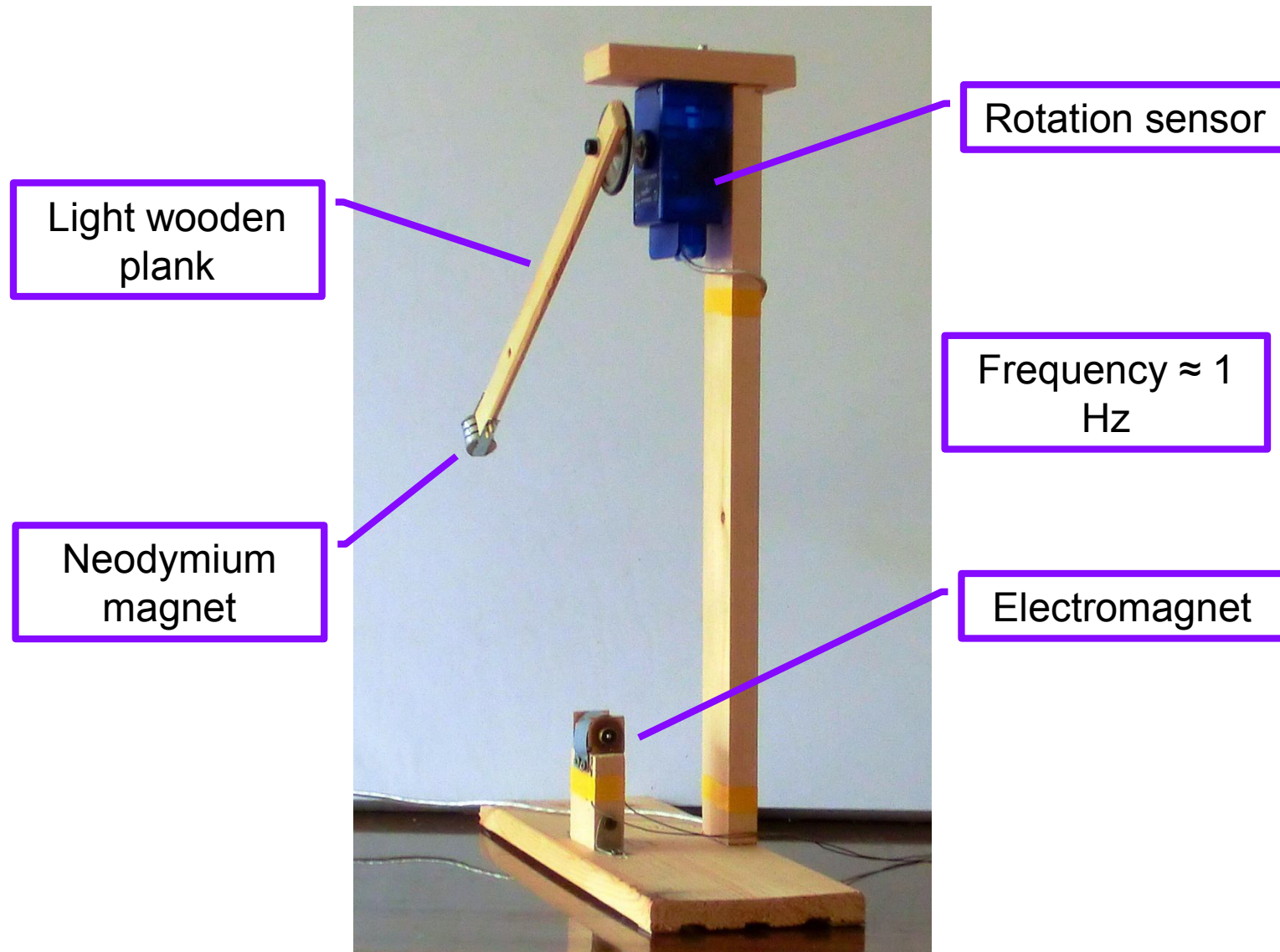


Make a light pendulum with a small magnet at the free end. An adjacent electromagnet connected to an AC power source of a much higher frequency than the natural frequency of the pendulum can lead to undamped oscillations with various amplitudes. Study and explain the phenomenon.

First observations

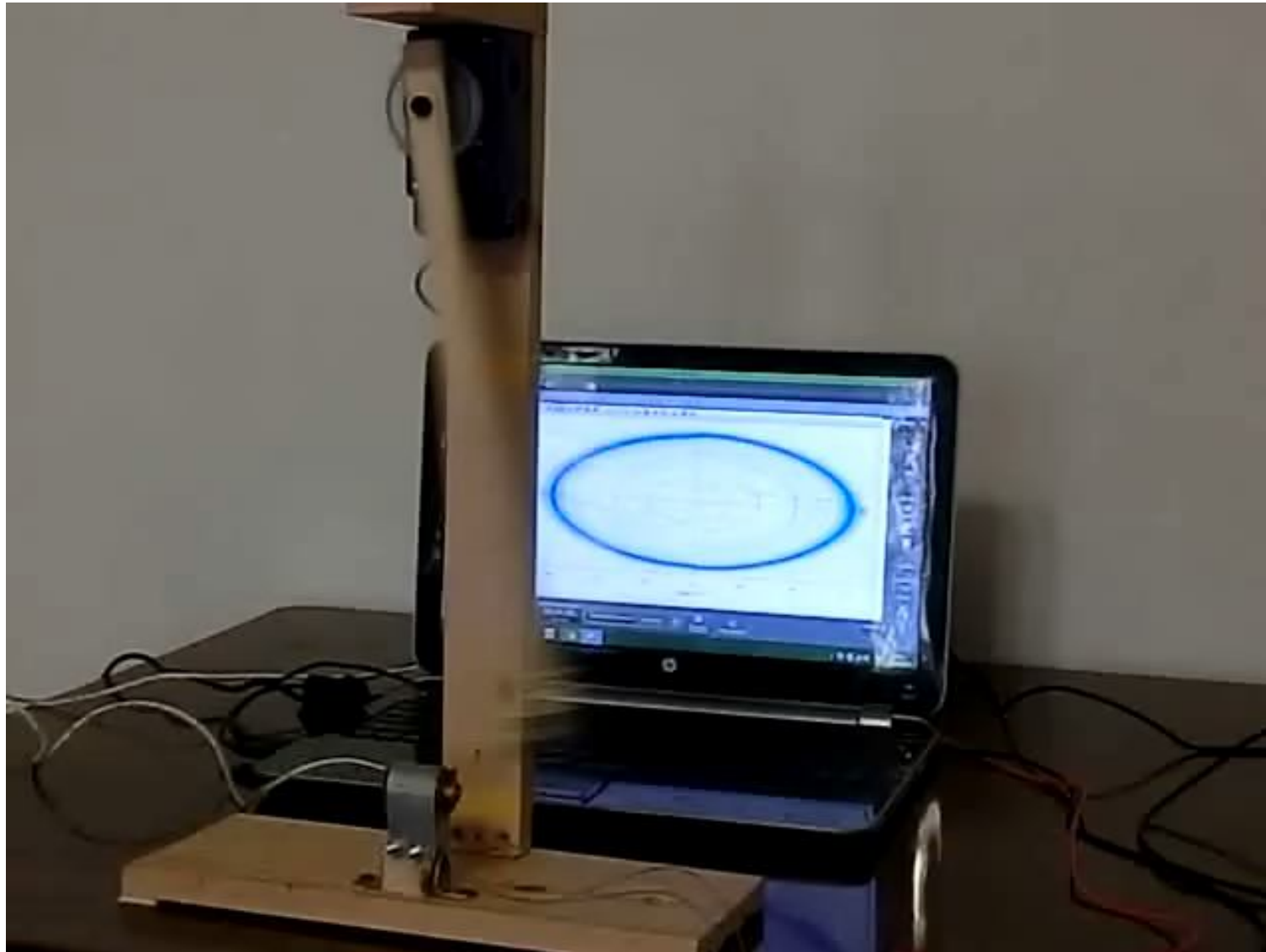
Experimental setup

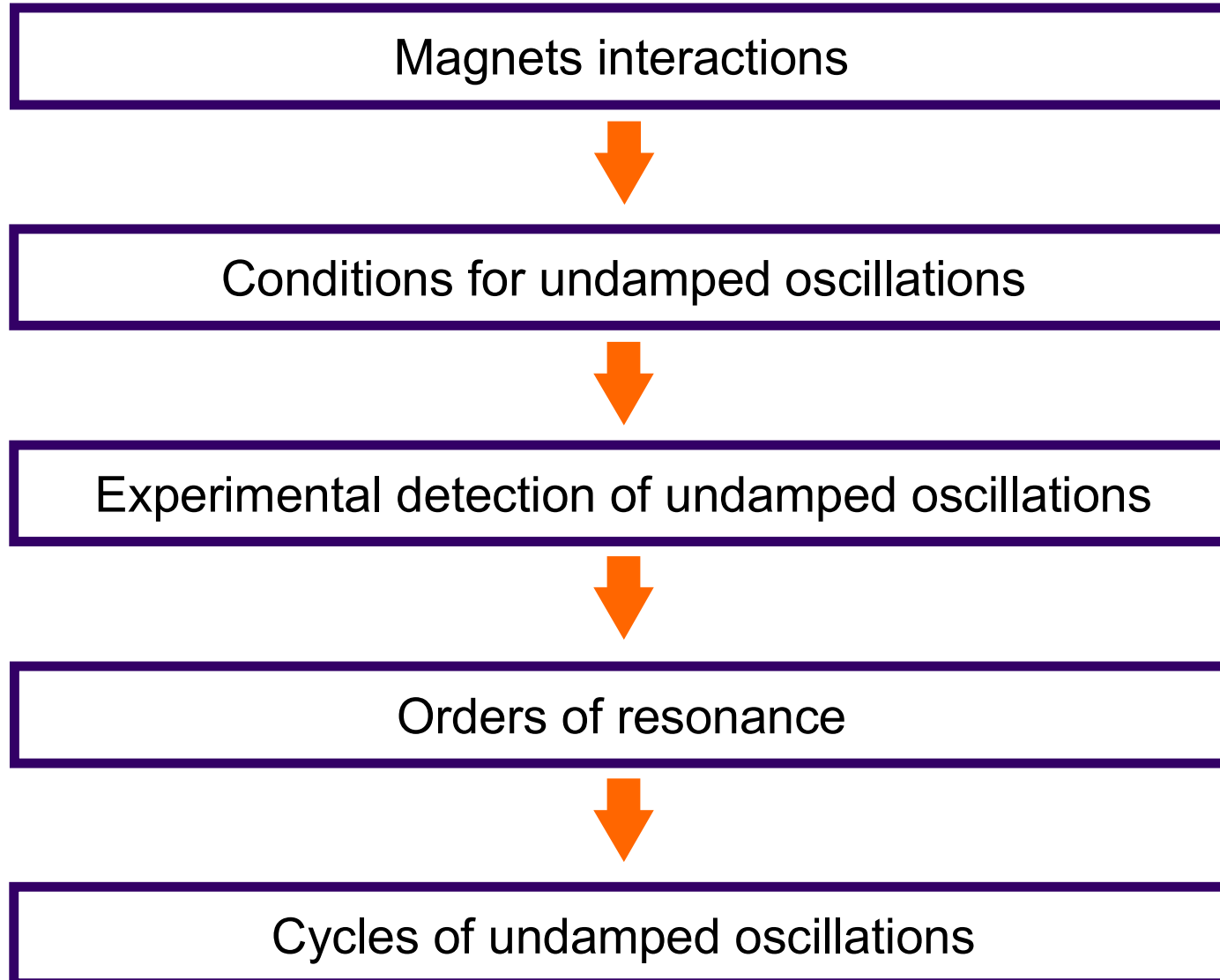
4



Undamped oscillations

5





Prediction of undamped oscillations



Computer simulation



Theoretical model



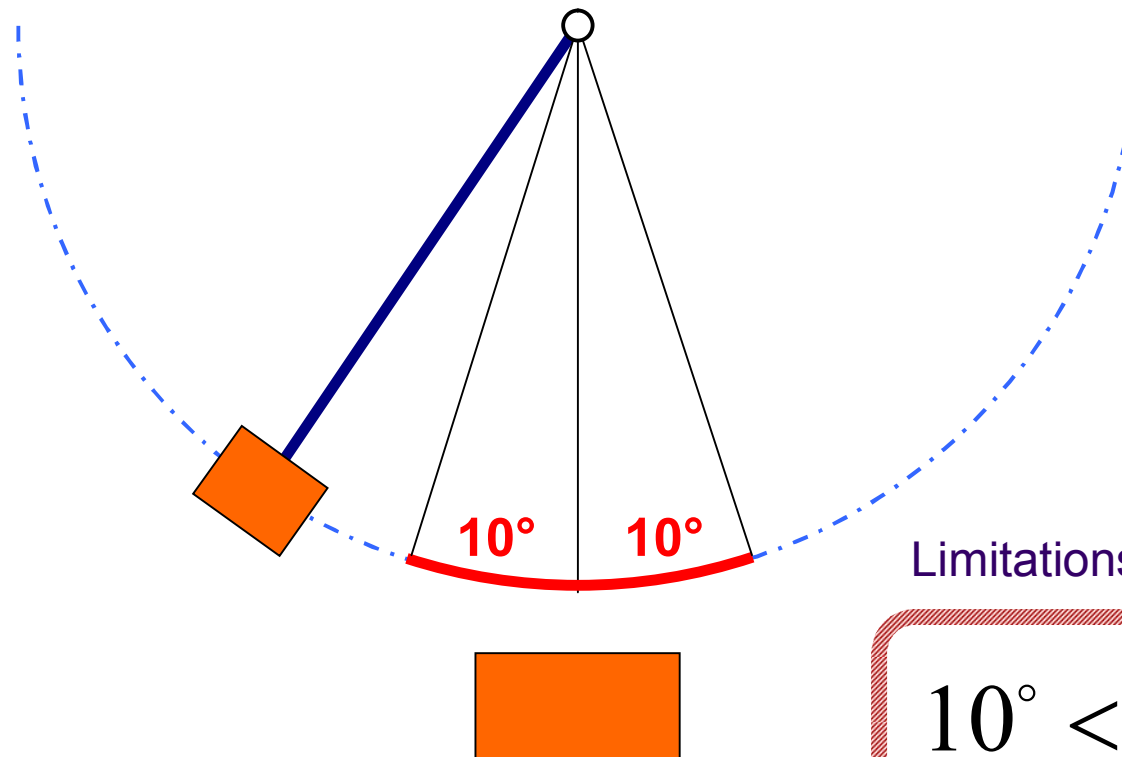
Poincare map

Qualitative explanation

Region of magnets' interaction

9

A magnet on the pendulum interacts with the electromagnet only flying over it. The span happens with almost constant velocity.



Limitations of amplitudes

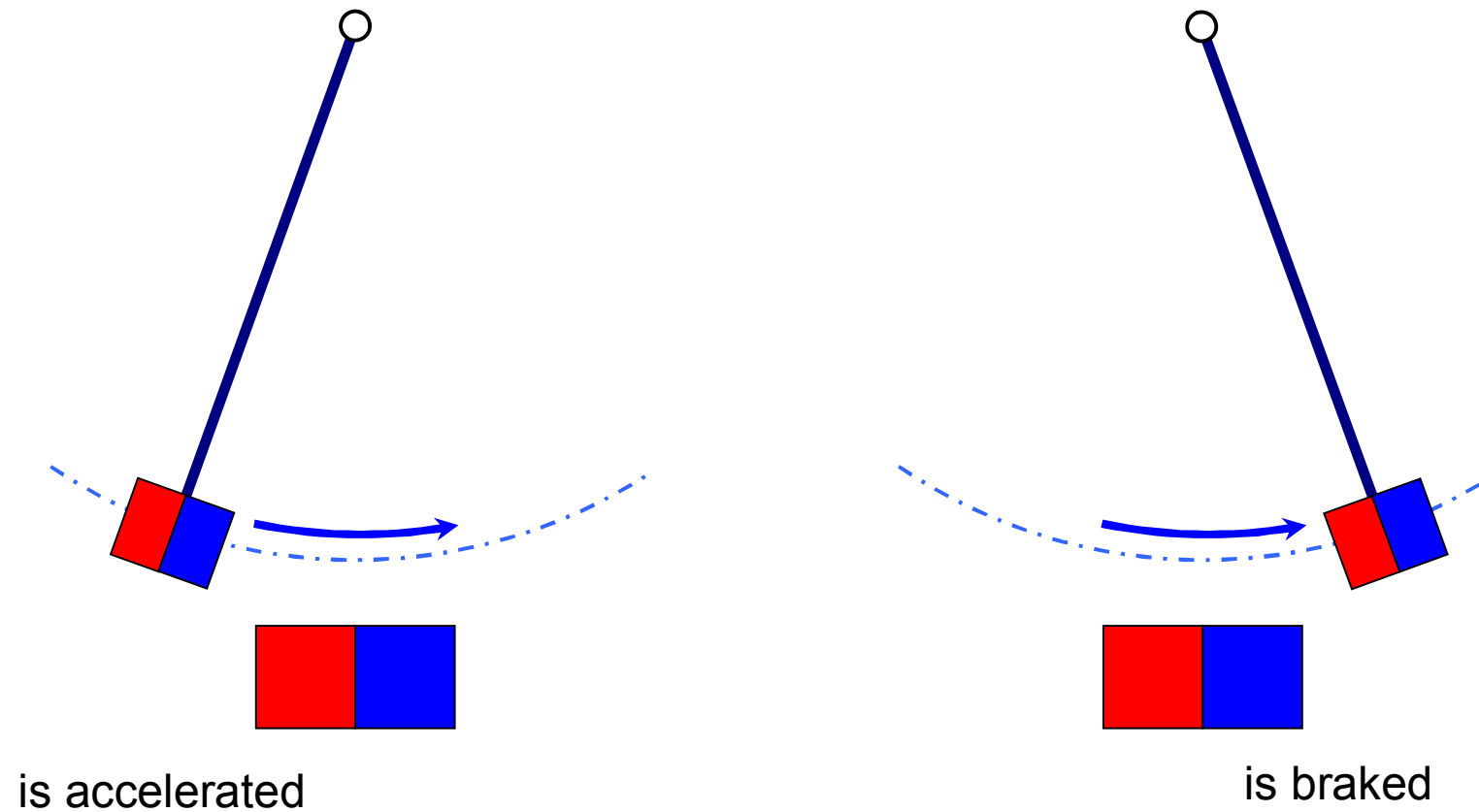
$$10^\circ < \alpha < 180^\circ$$

$$v = \sqrt{2gl} \approx 2.5 \text{ m/s}$$

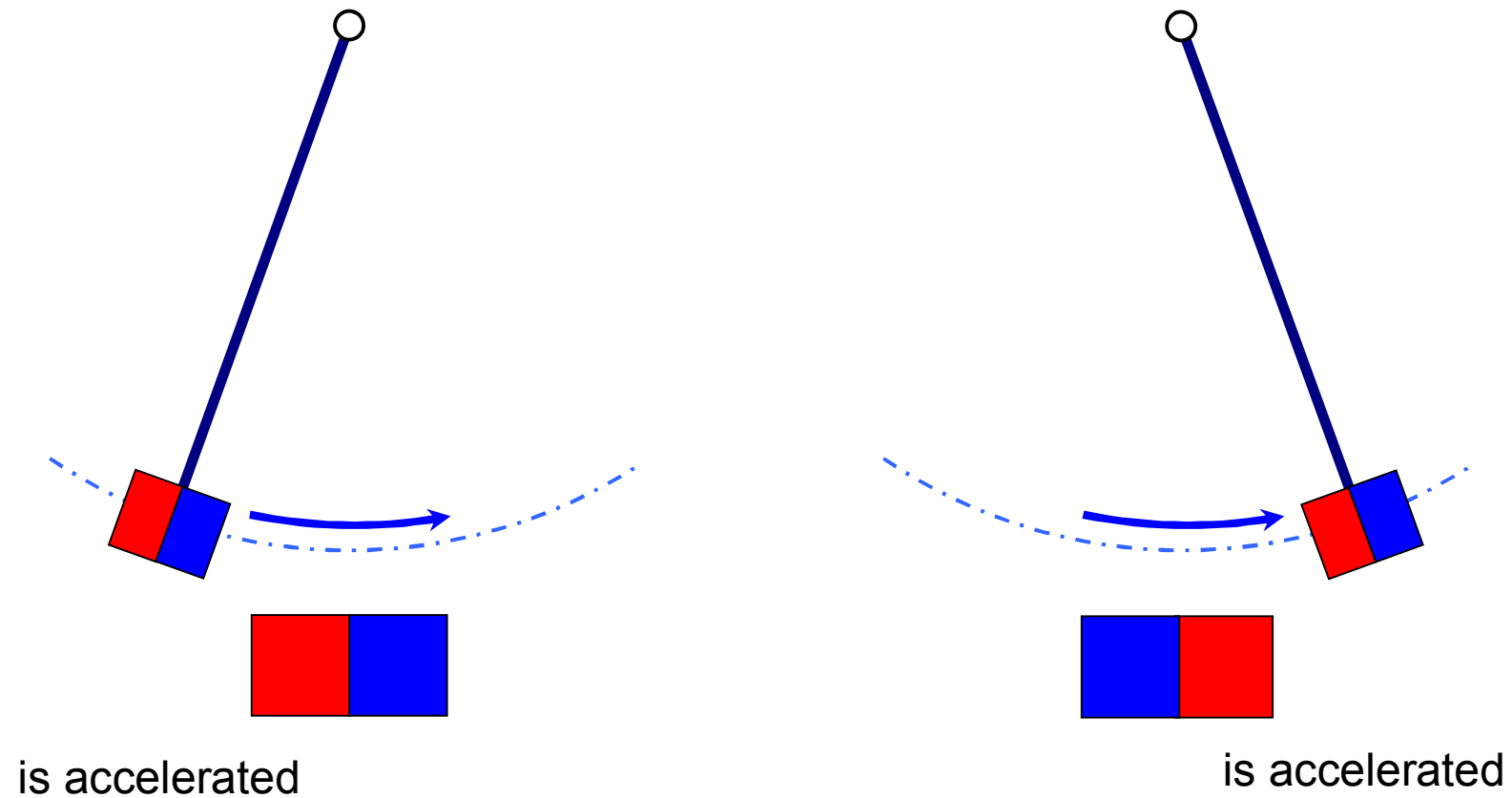
$$\tau = \frac{2d}{v} \approx 0.04 \text{ s}$$

$$f = \frac{1}{\tau} \approx 25 \text{ Hz}$$

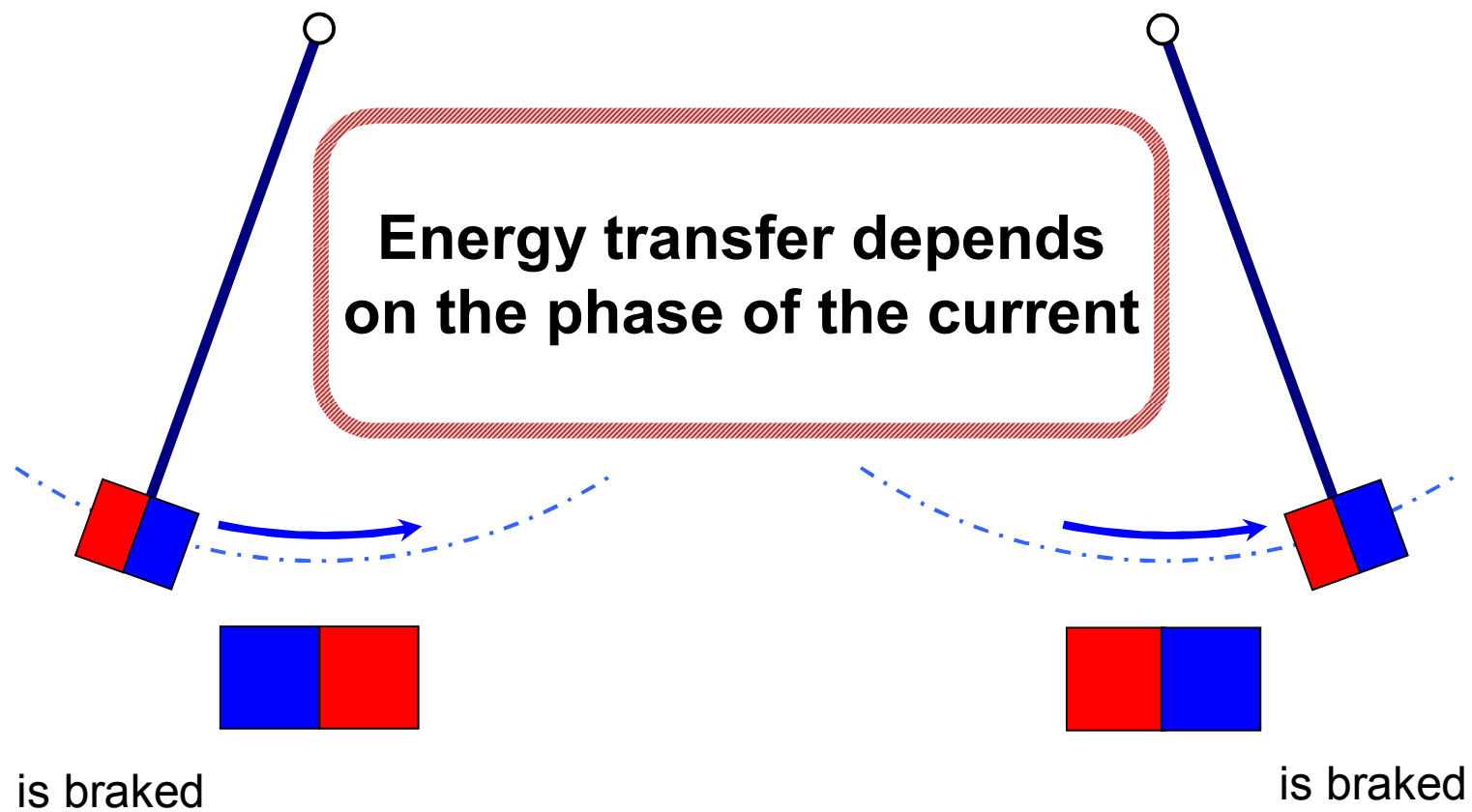
If the current in solenoid is direct, the change in energy of the pendulum is equal to zero.



The energy of the pendulum can either increase...



...and decrease.

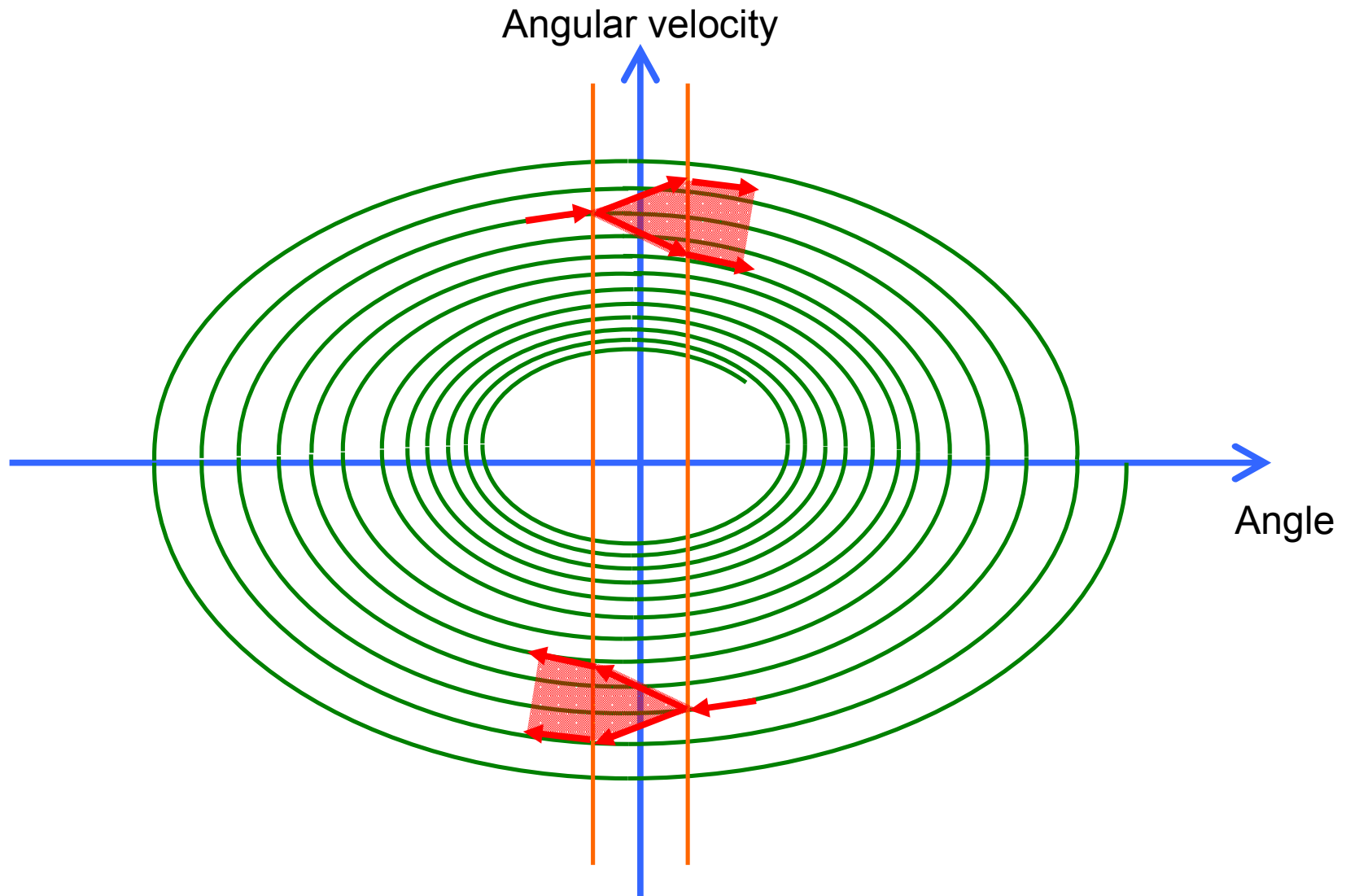


We denote the time when the pendulum passes its lowest point as $t = 0$.

Current oscillations in the coil:

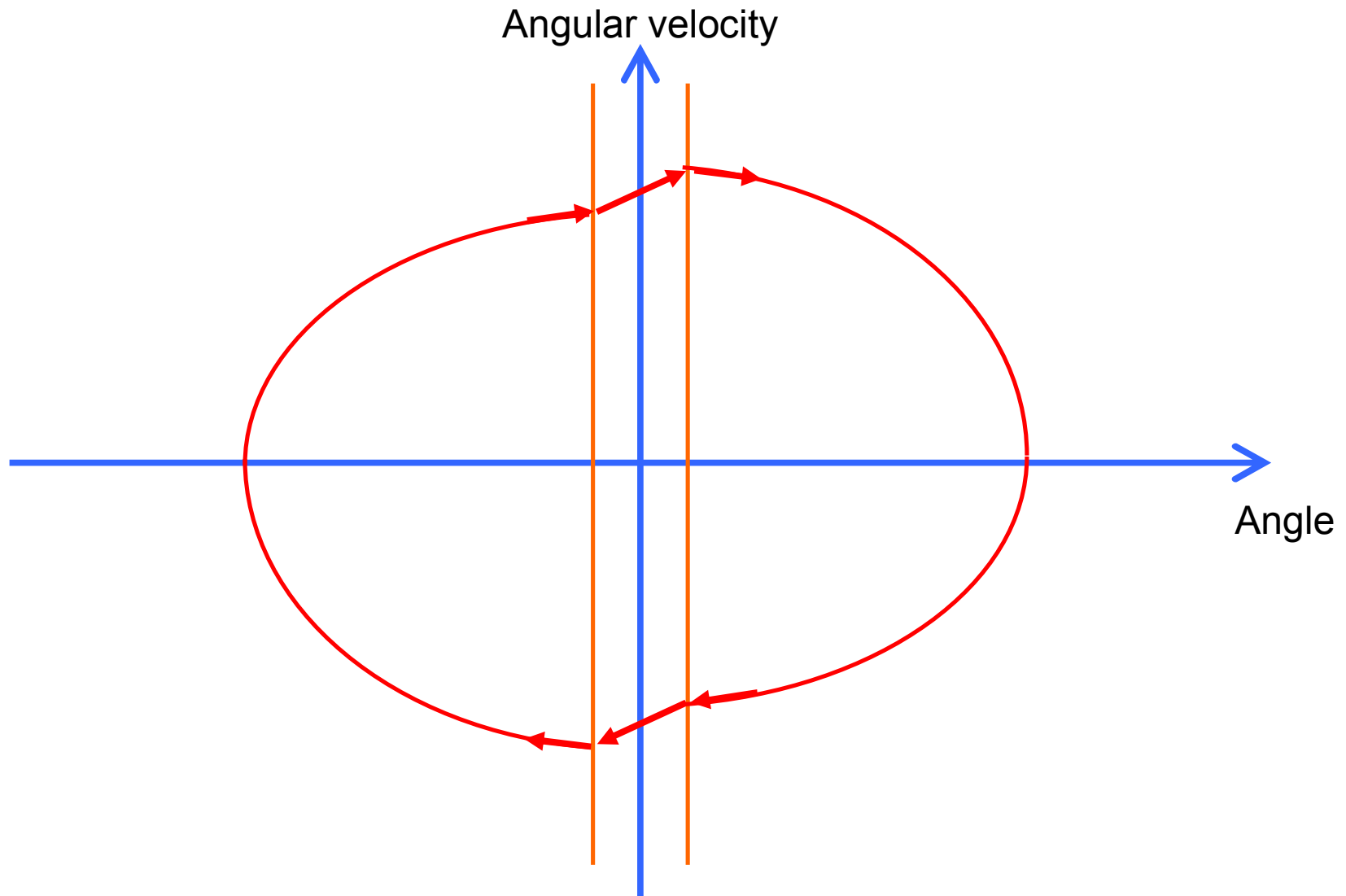
$$I(t) = I_0 \sin(\omega t + \varphi)$$

$$I(t) = \underbrace{(I_0 \cos \varphi) \cdot \sin \omega t}_{\text{Energy transfer}} + \underbrace{(I_0 \sin \varphi) \cdot \cos \omega t}_{\text{No energy transfer}}$$



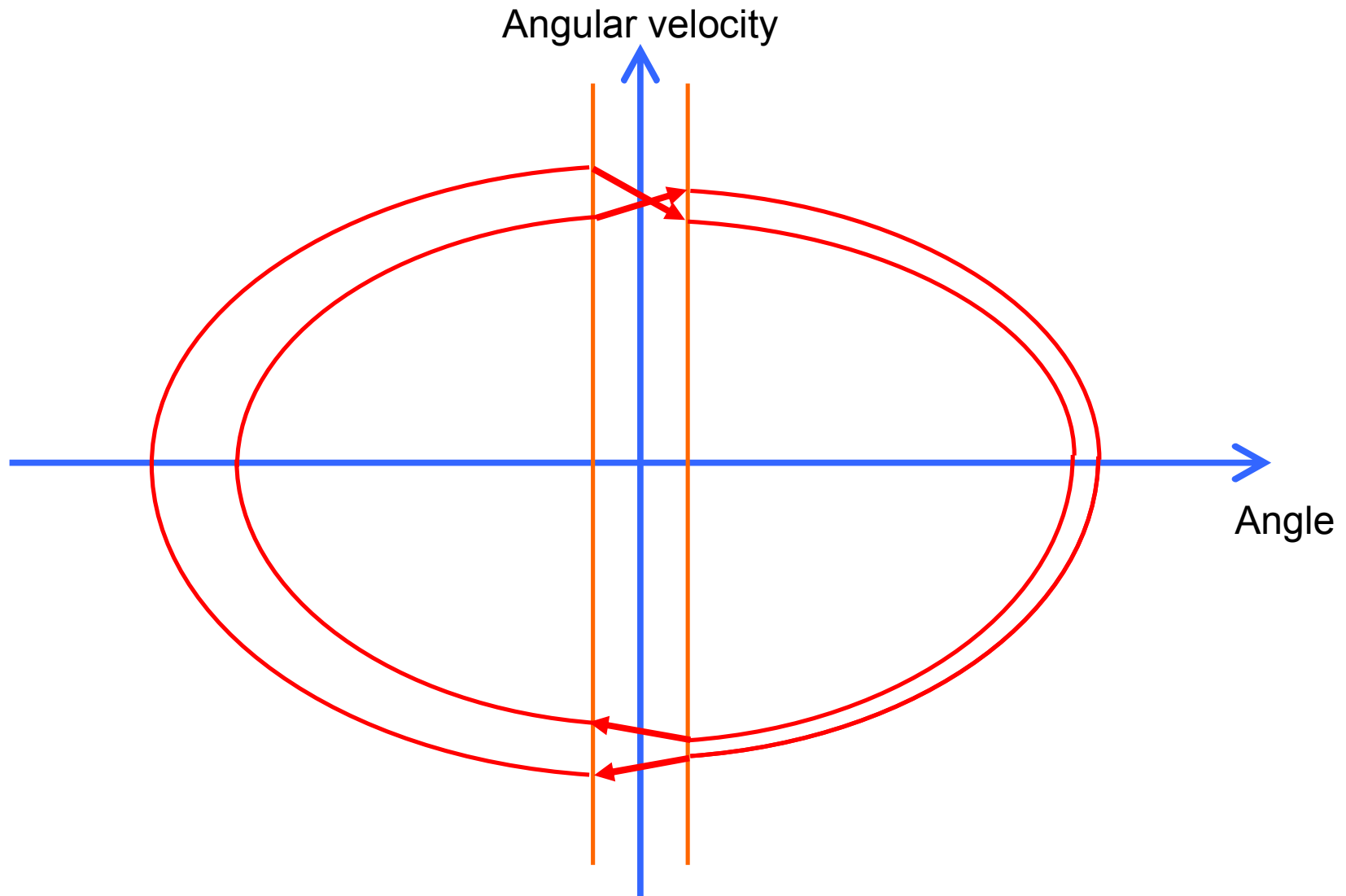
Undamped oscillations (1 period)

16

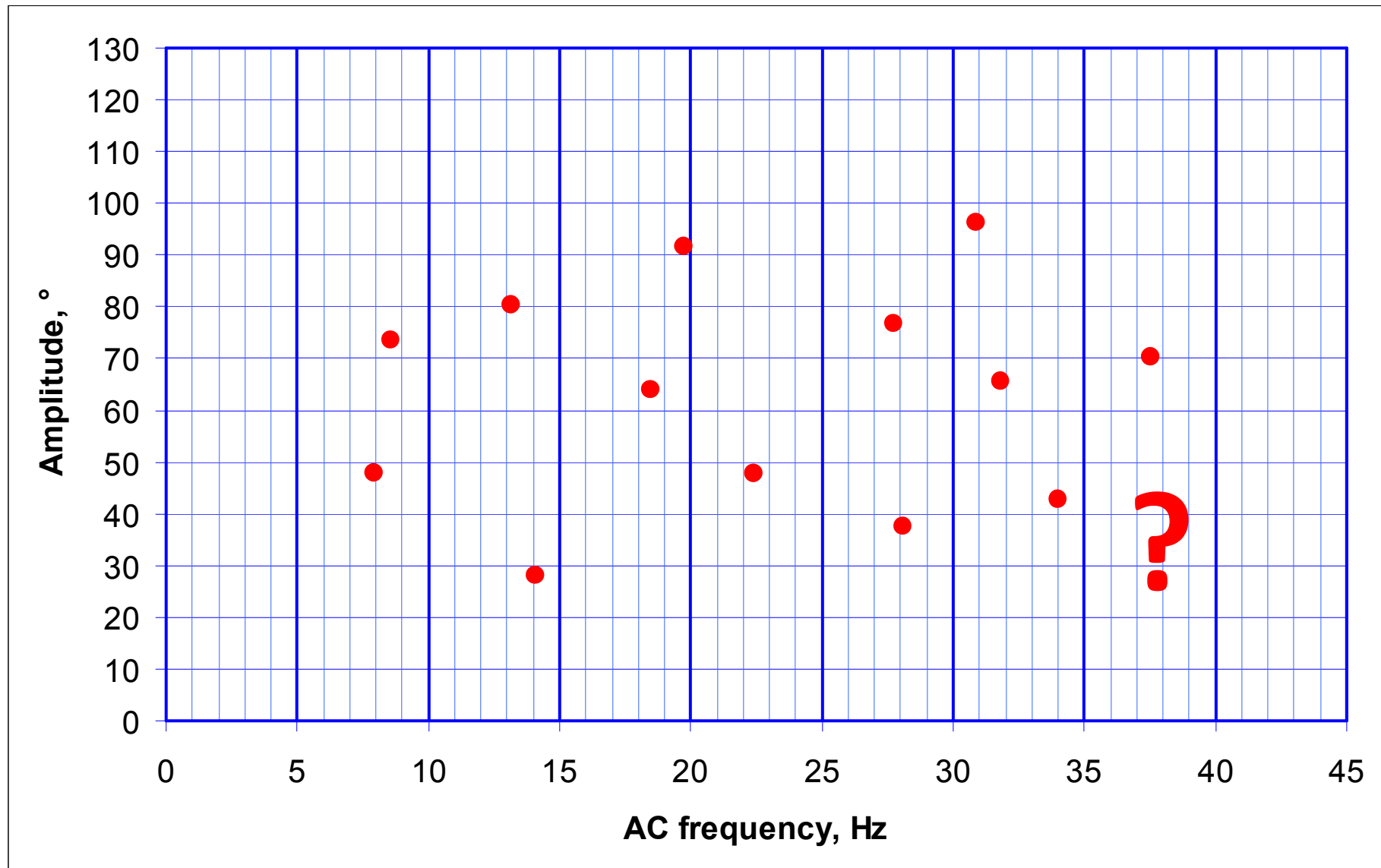


Undamped oscillations (2 periods)

17

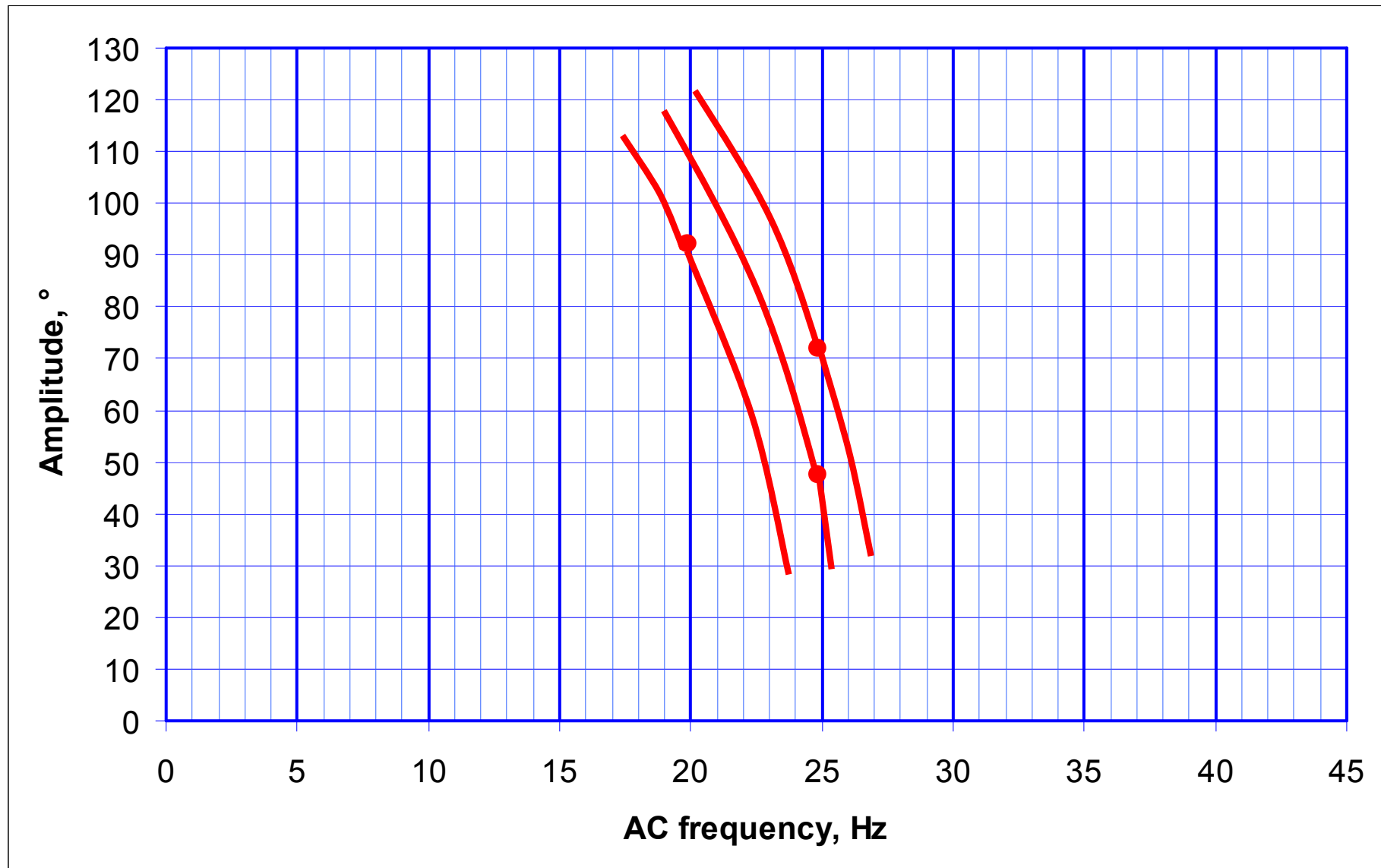


Experiments



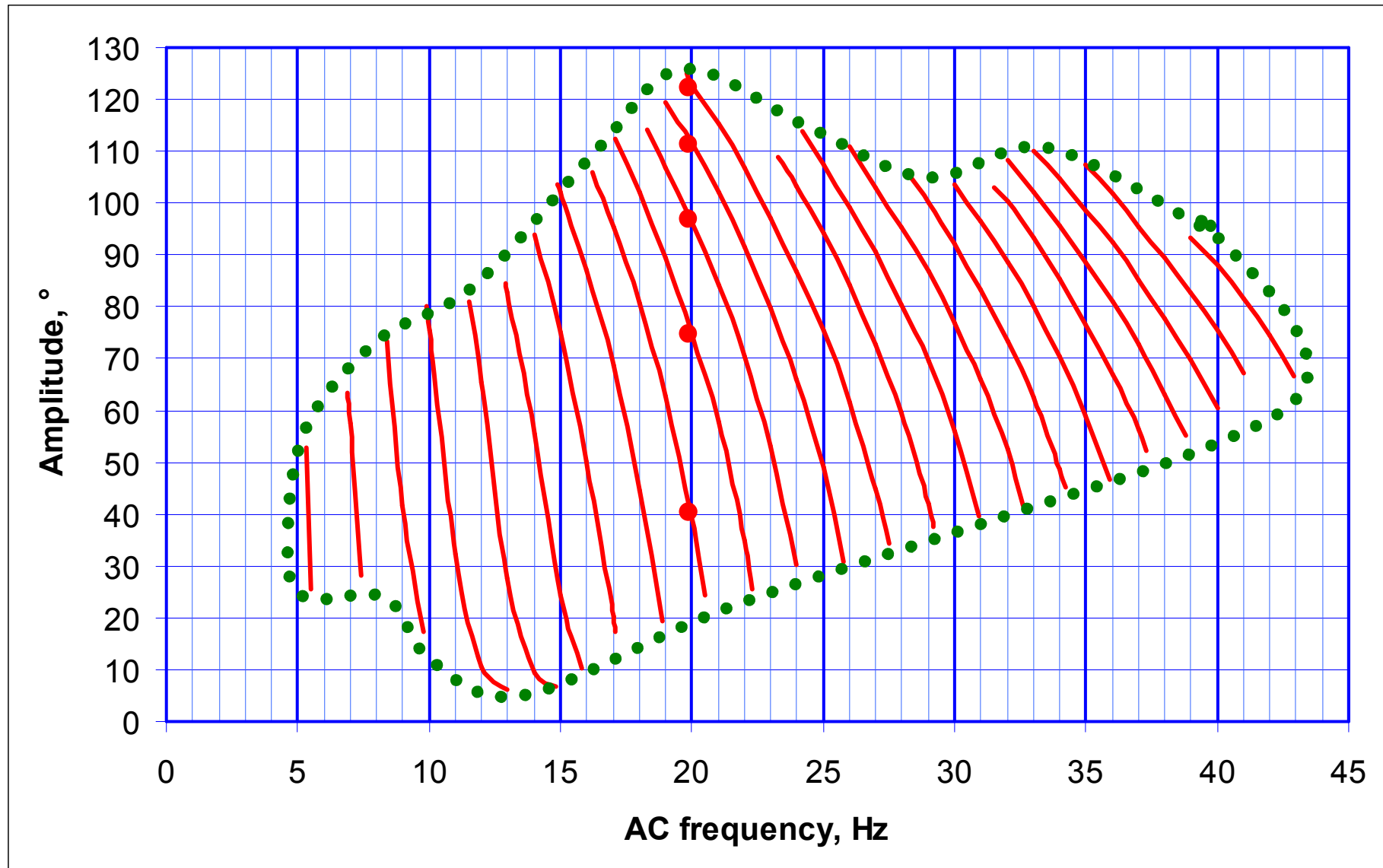
Experiment #1

20



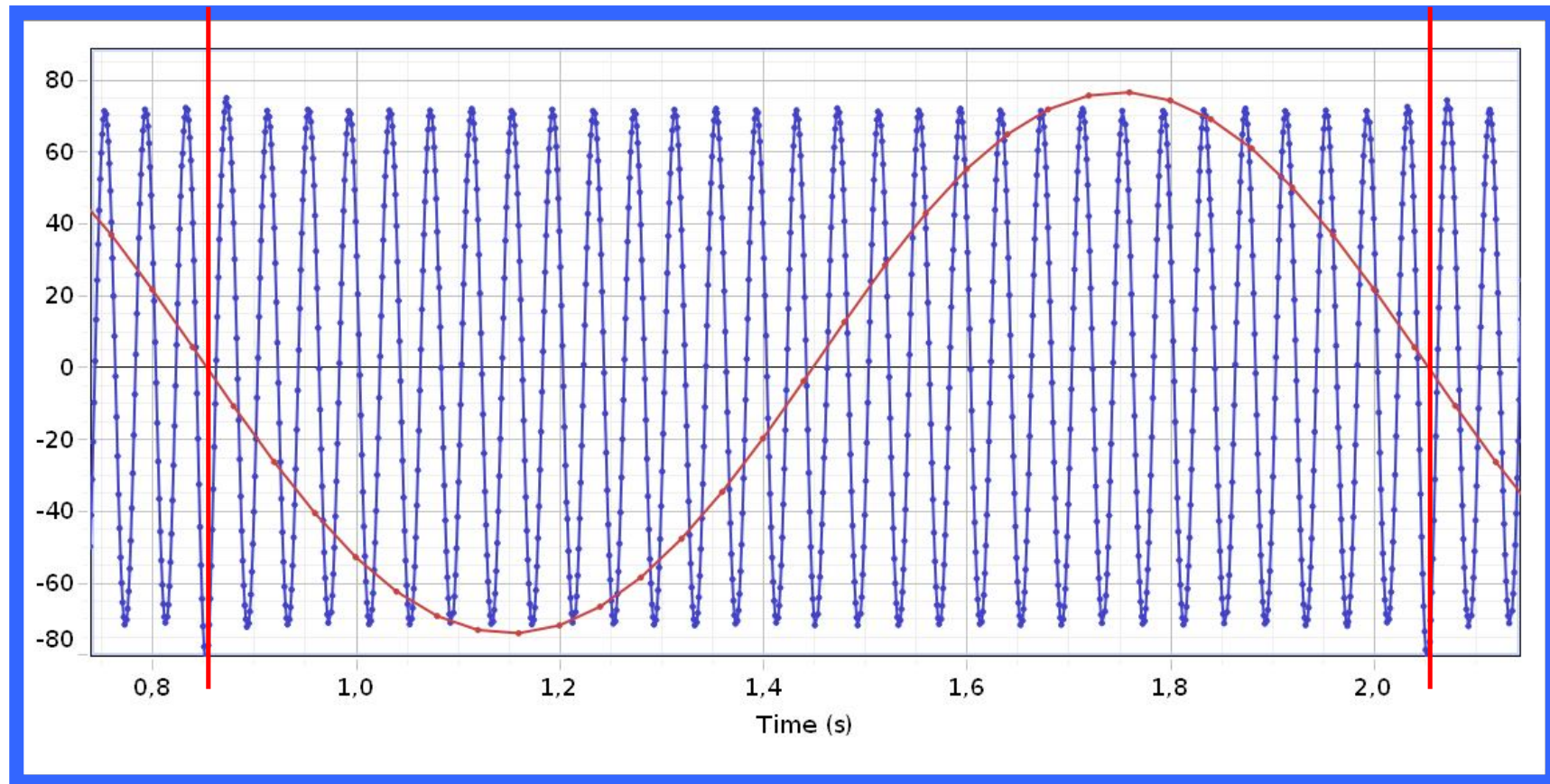
Undamped oscillations

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Order of resonance

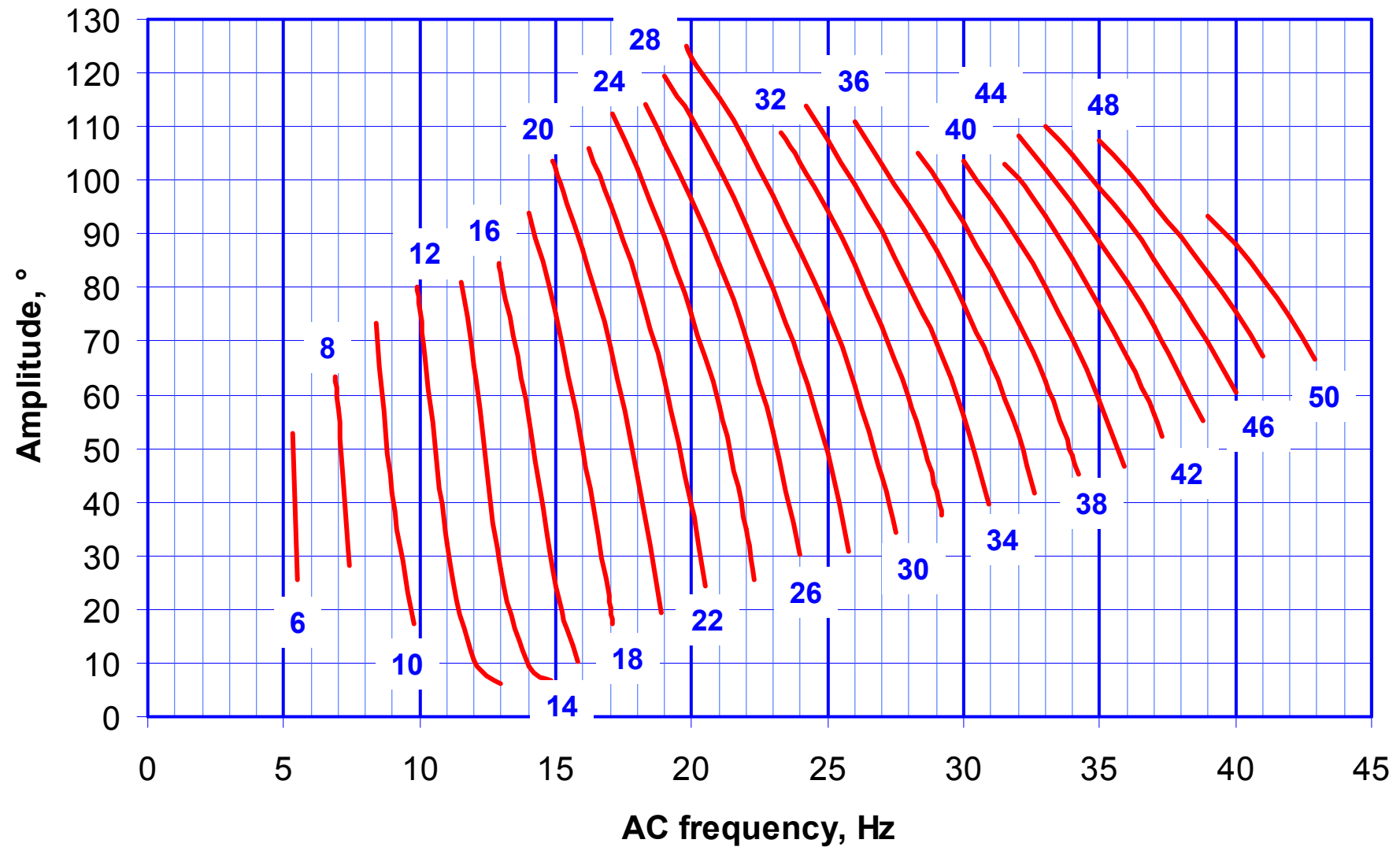
22



In this test 30 current oscillations occur during
1 pendulum oscillation. Order of resonance = 30.

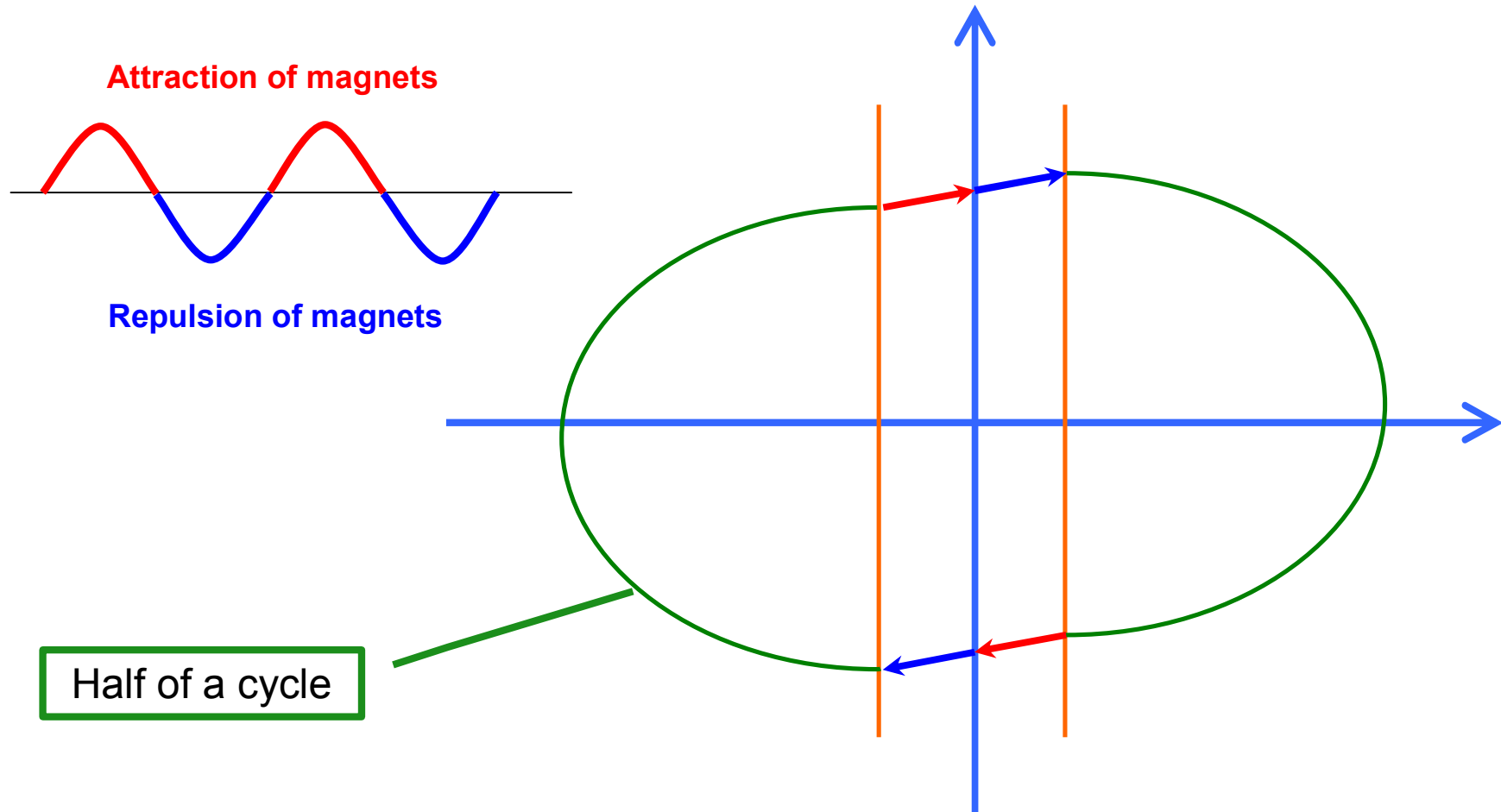
Orders of resonance

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Why all orders are even?

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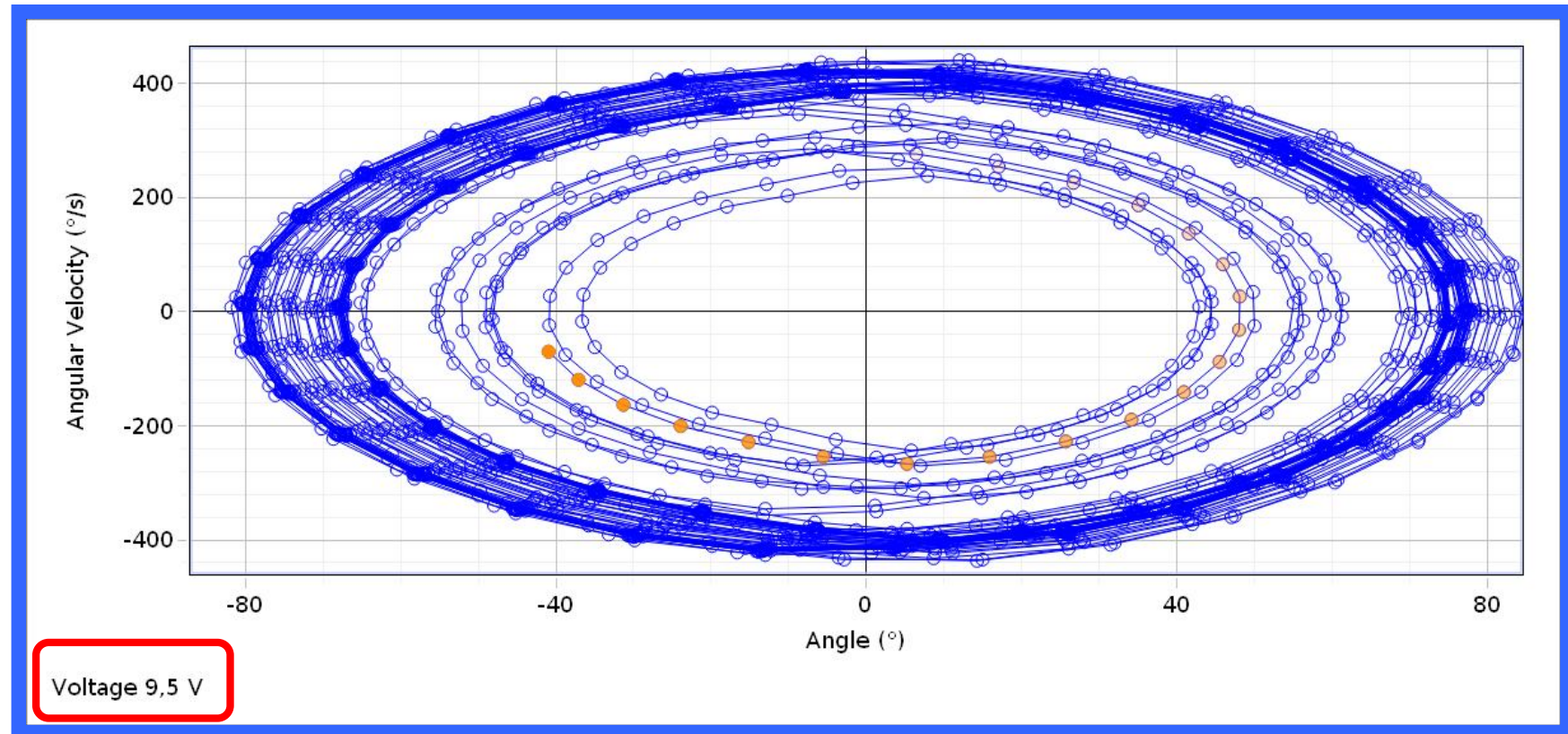
$$N=2+2\cdot k$$

Cycles of undamped oscillations

- AC frequency is fixed (25 Hz in the main series).
- Some undamped oscillation of kicked pendulum is found (resonance order = 30, amplitude $\approx 80^\circ$ in the main series).
- Slowly shifting the voltage we monitor changes of the fine structure of the phase portrait.

Period doubling

27



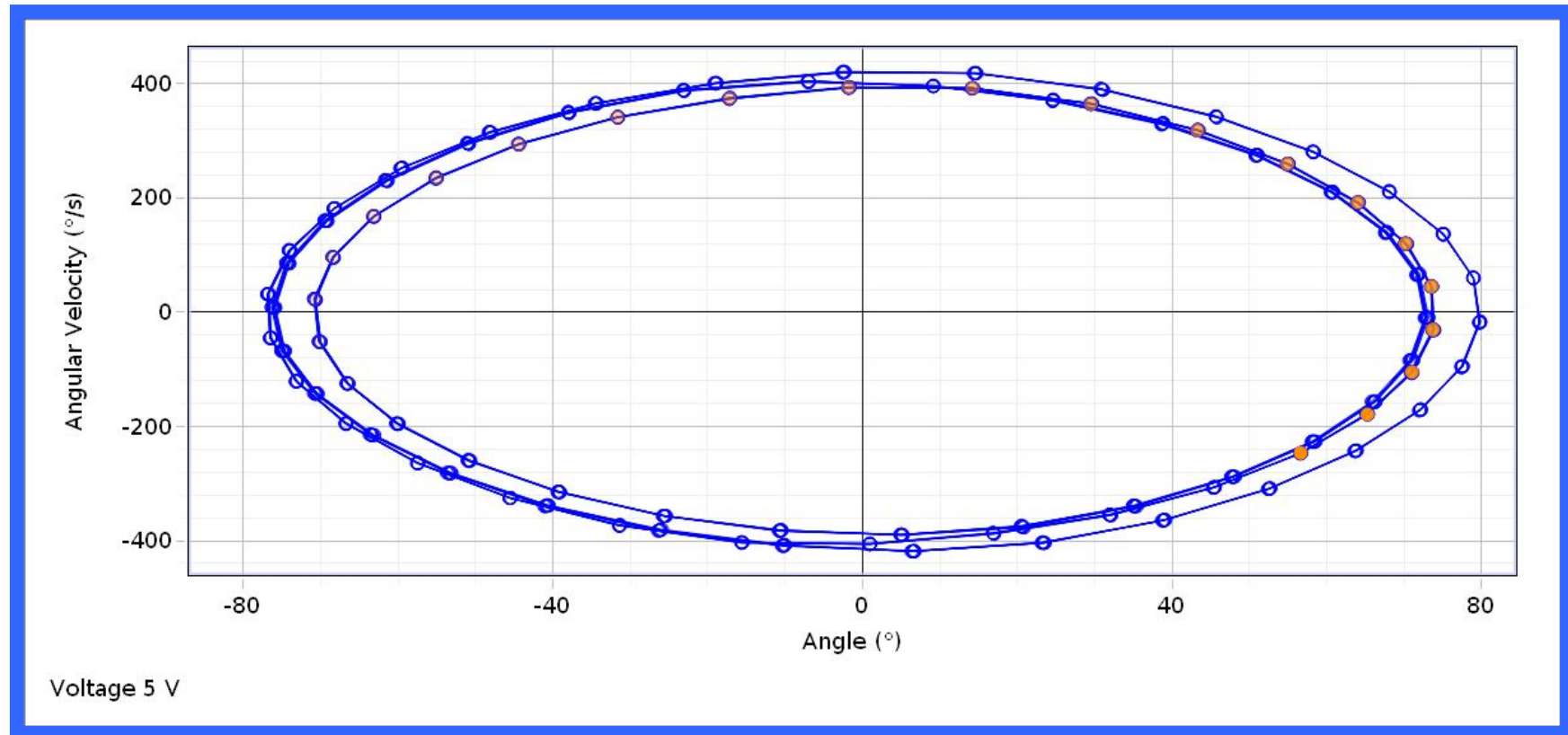
Appearance of the cycle = **1.7 V**

Period doubling = **6.9 V**

Destruction of the cycle = **9.5 V**

3-periodic cycle

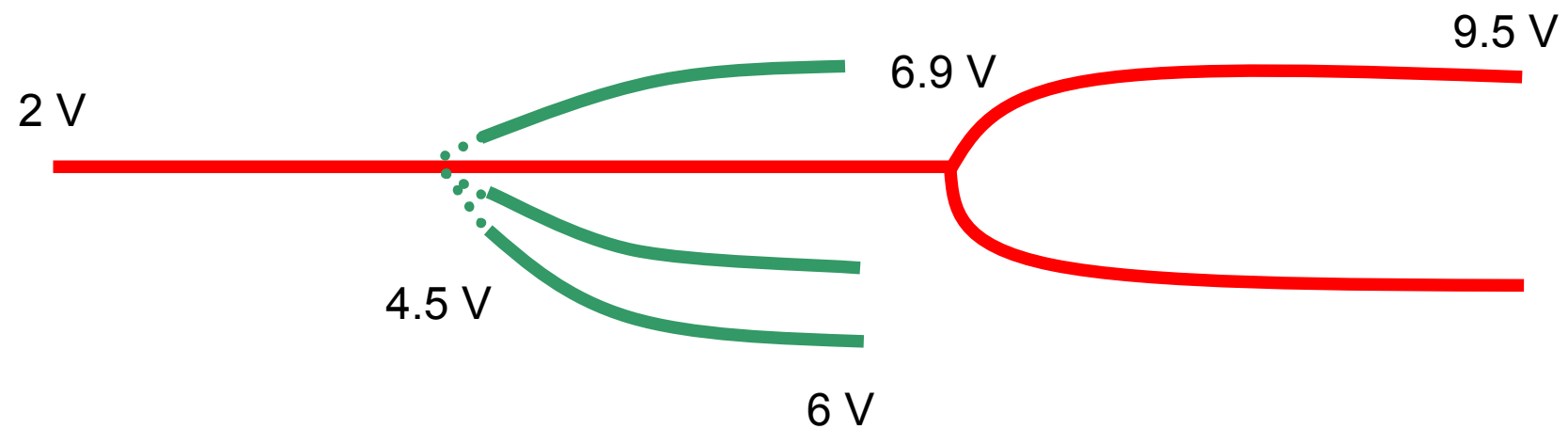
28



This 3-periodic cycle exists
in the range from **4.5 V** to **6.0 V**

The fine structure of the resonance

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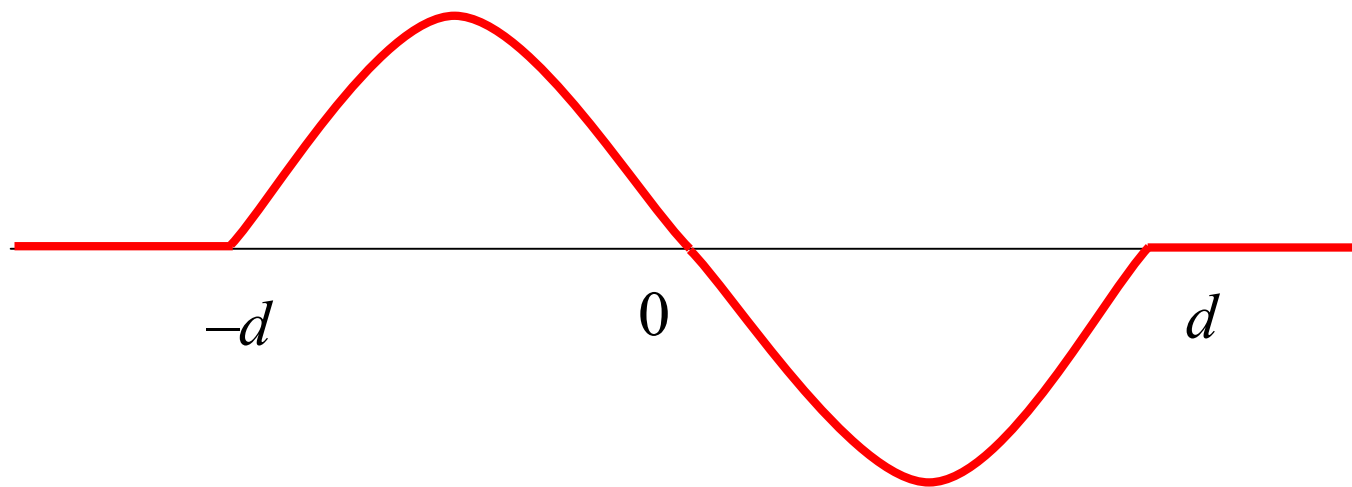


- **At low currents:** kicks are too weak to compensate the energy dissipation.
- **At high currents:** phase adjustment can't exactly compensate too strong too strong kicks. This eventually leads to destruction of the cycle.

Computer simulation

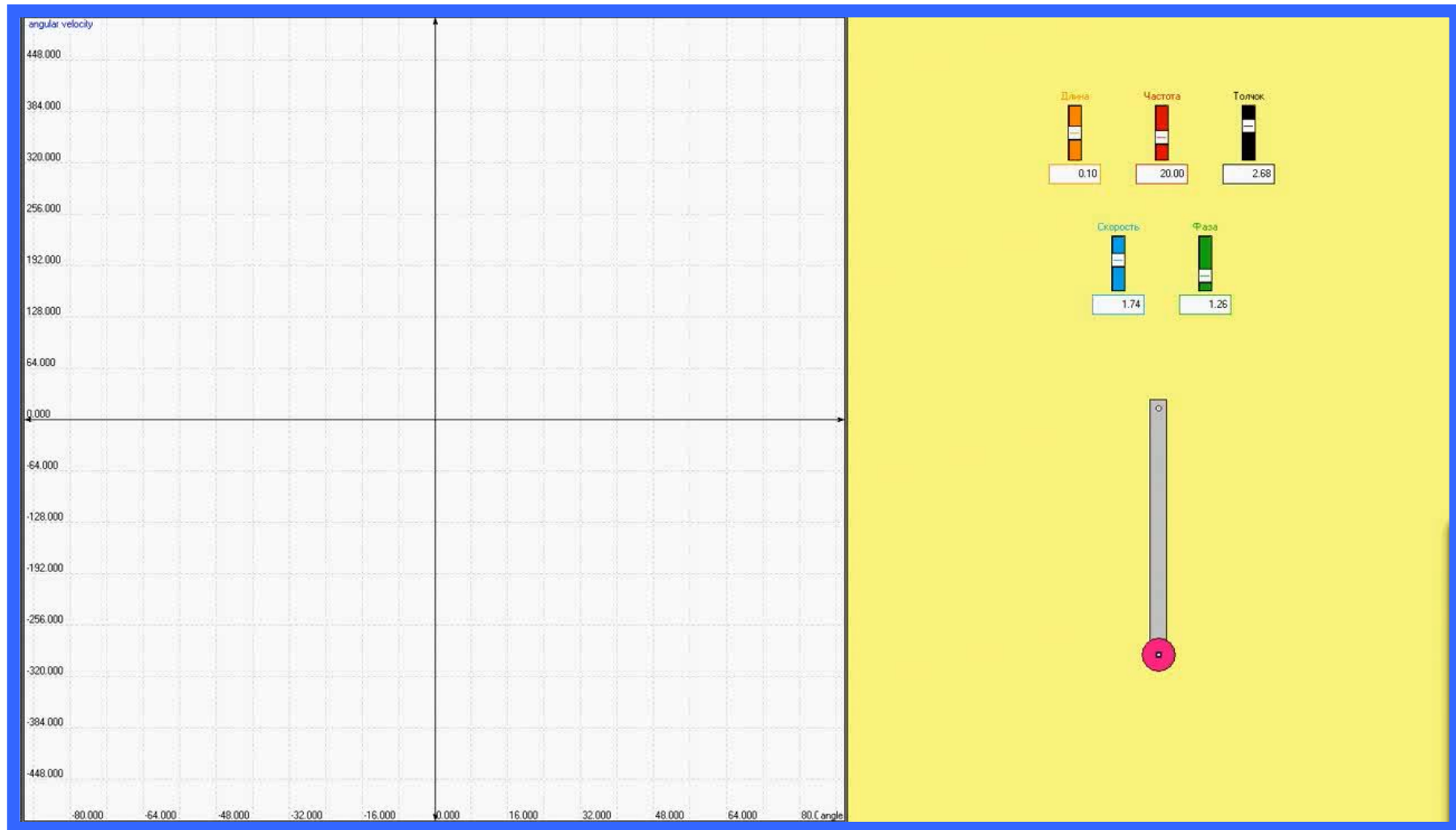
- There are limitations on the voltage and current in power supply.
- Fine effects are difficult to observe because of vibrations and external disturbances.

$$F(x) = F_0 \sin\left(\frac{\pi x}{d}\right)$$



Computer simulation

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- We do not know which initial start conditions lead to the establishment of undamped oscillations.
- In each test, we need to wait 10 minutes to establish the stationary regime.
- To conduct all the necessary measurements it will take us more then decade of continuous work!

Poincare map

The pendulum interacts with an
electromagnet in the narrow angle range.
This allows to relate values (v, φ) on two
consecutive spans.

$$\begin{pmatrix} v_n \\ \varphi_n \end{pmatrix} \quad \boxed{\begin{pmatrix} v_n \\ \varphi_n \end{pmatrix} \xrightarrow{\text{Theoretical model}} \begin{pmatrix} v_{n+1} \\ \varphi_{n+1} \end{pmatrix}} \quad \begin{pmatrix} v_{n+1} \\ \varphi_{n+1} \end{pmatrix}$$

The electromagnet is switched off:

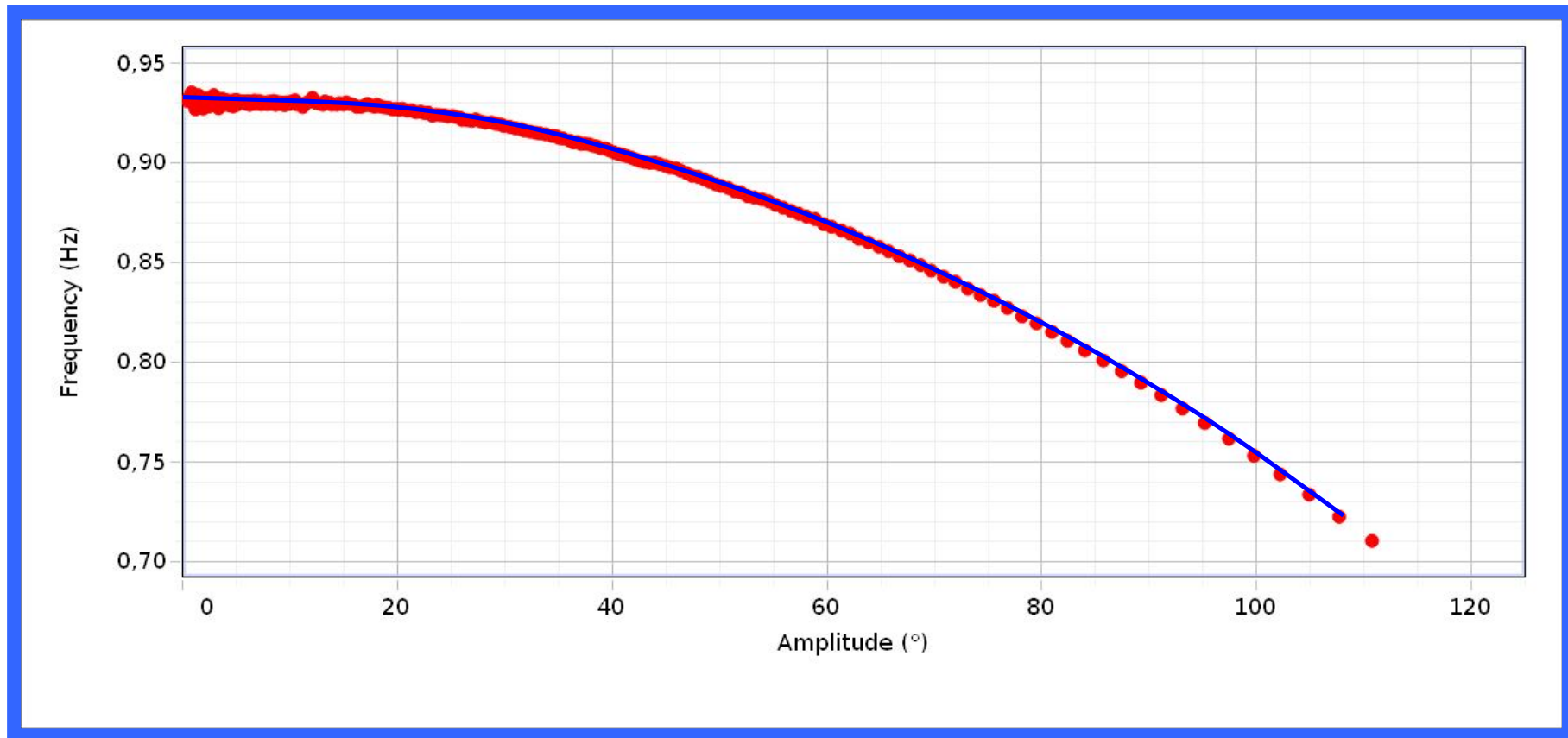
$$v_{n+1} = (1 - \varepsilon) \cdot v_n$$

The electromagnet is switched on:

$$v_{n+1} = (1 - \varepsilon) \cdot v_n + \frac{F_0}{m\omega} \cdot \frac{2\pi \left(\frac{\omega d}{v_n} \right) \cdot \sin \left(\frac{\omega d}{v_n} \right)}{\left(\pi^2 - \left(\frac{\omega d}{v_n} \right)^2 \right)} \cdot \cos \varphi_n$$

Oscillations with large amplitude

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$$f(\alpha) = f_0 \cdot \frac{\pi / 2}{\int_0^\alpha \sqrt{2(\cos \theta - \cos \alpha)} d\theta}$$

Change in phase for half-period

40

Dependence of the oscillation period on the velocity.

$$T(v) = T_0 \cdot \left\{ 1 + \left(\frac{1}{2} \right)^2 \left(\frac{v^2}{2U^2} \right) + \left(\frac{1 \cdot 3}{2 \cdot 4} \right)^2 \left(\frac{v^2}{2U^2} \right)^2 + \left(\frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6} \right)^2 \left(\frac{v^2}{2U^2} \right)^3 + \dots \right\}$$

Our calculation takes into account
first six expansion terms.

$$U^2 = 2gl$$

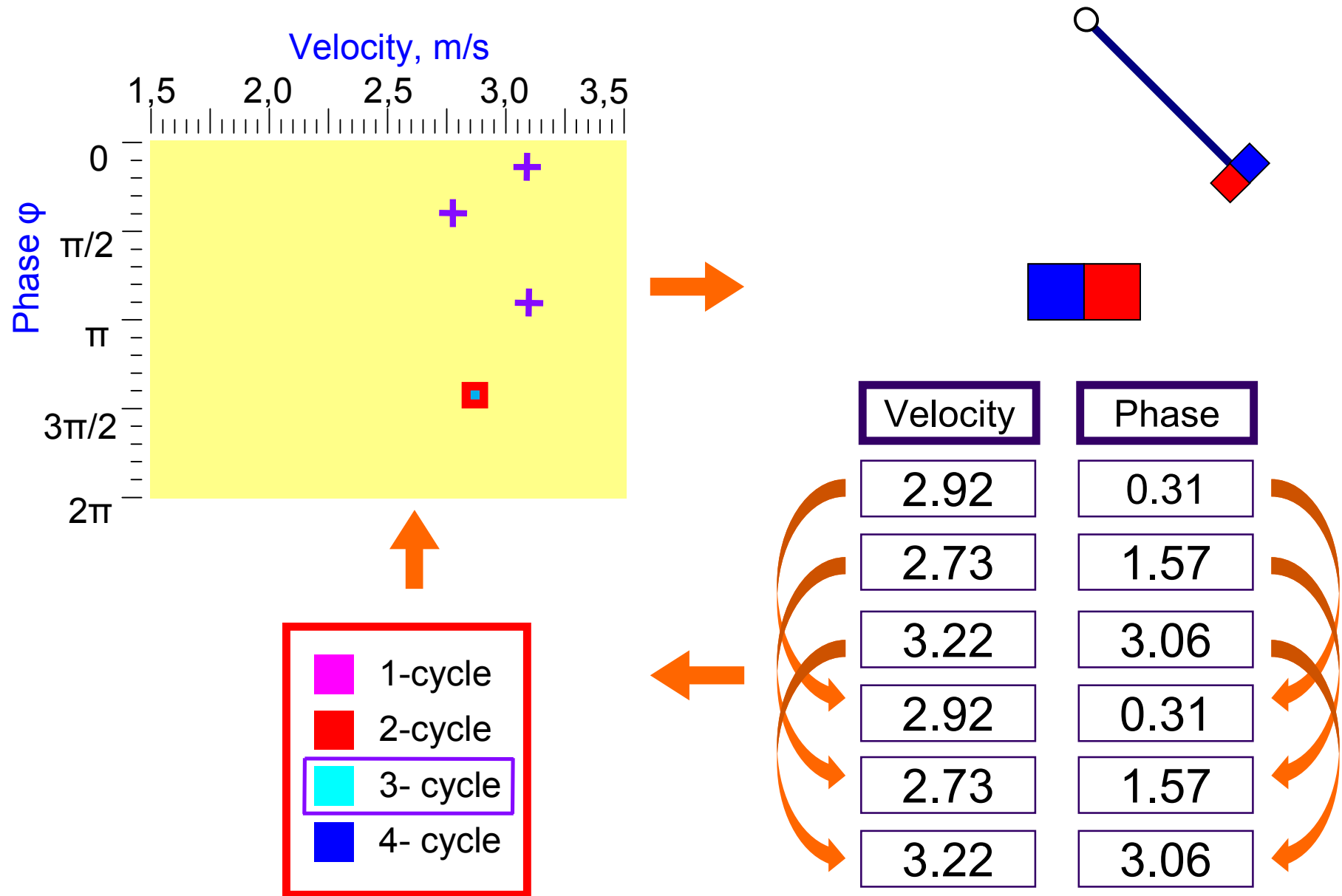
$$\varphi_{n+1} = \varphi_n + \omega \cdot \frac{T(v_{n+1})}{2}$$

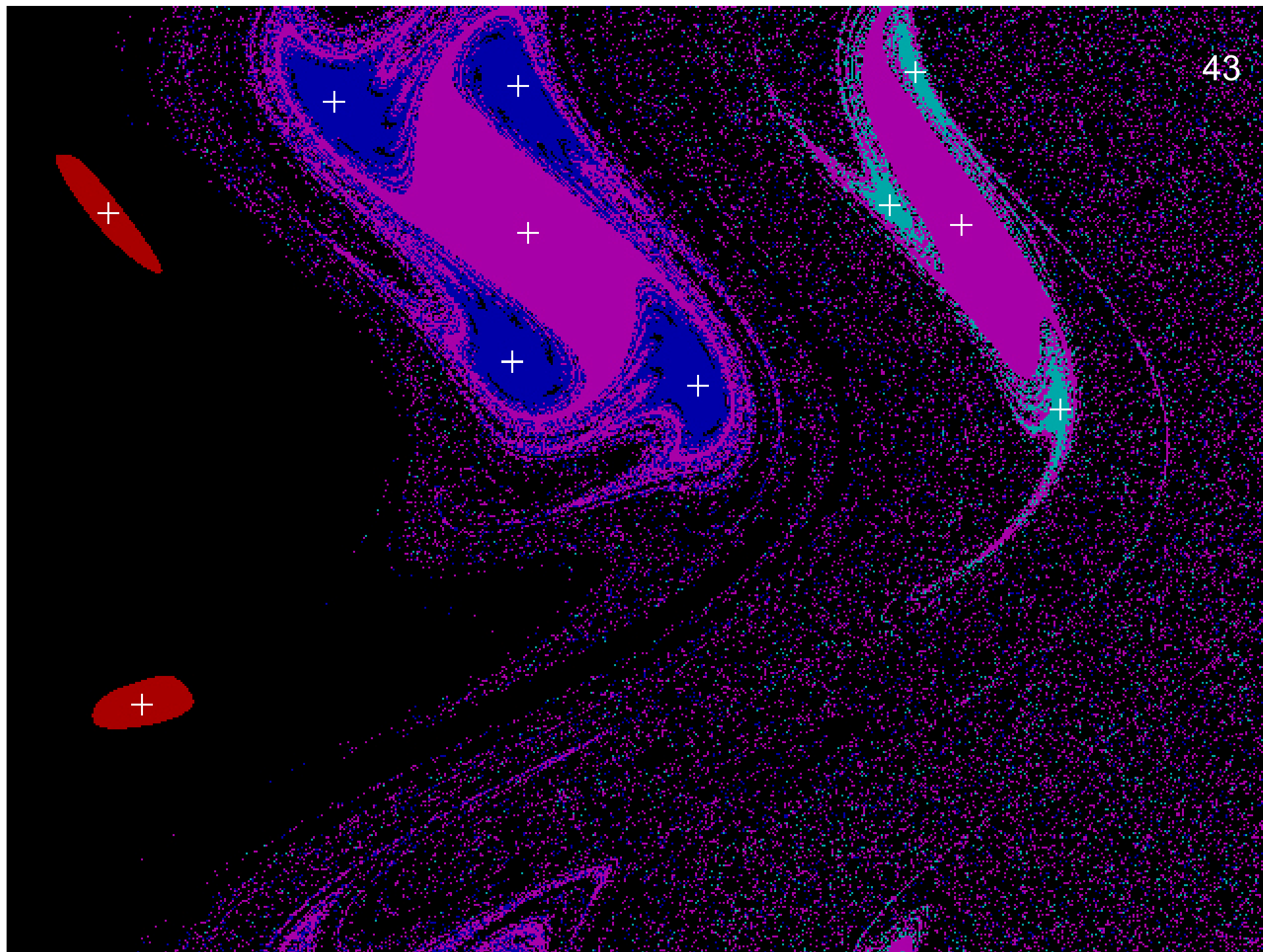
$$v_{n+1} = (1 - \varepsilon) \cdot v_n + \frac{F_0}{m\omega} \cdot \frac{2\pi \left(\frac{\omega d}{v_n} \right) \cdot \sin \left(\frac{\omega d}{v_n} \right)}{\left(\pi^2 - \left(\frac{\omega d}{v_n} \right)^2 \right)} \cdot \cos \varphi_n$$

$$\varphi_{n+1} = \varphi_n + \omega \cdot \frac{T(v_{n+1})}{2}$$

Poincare mapping

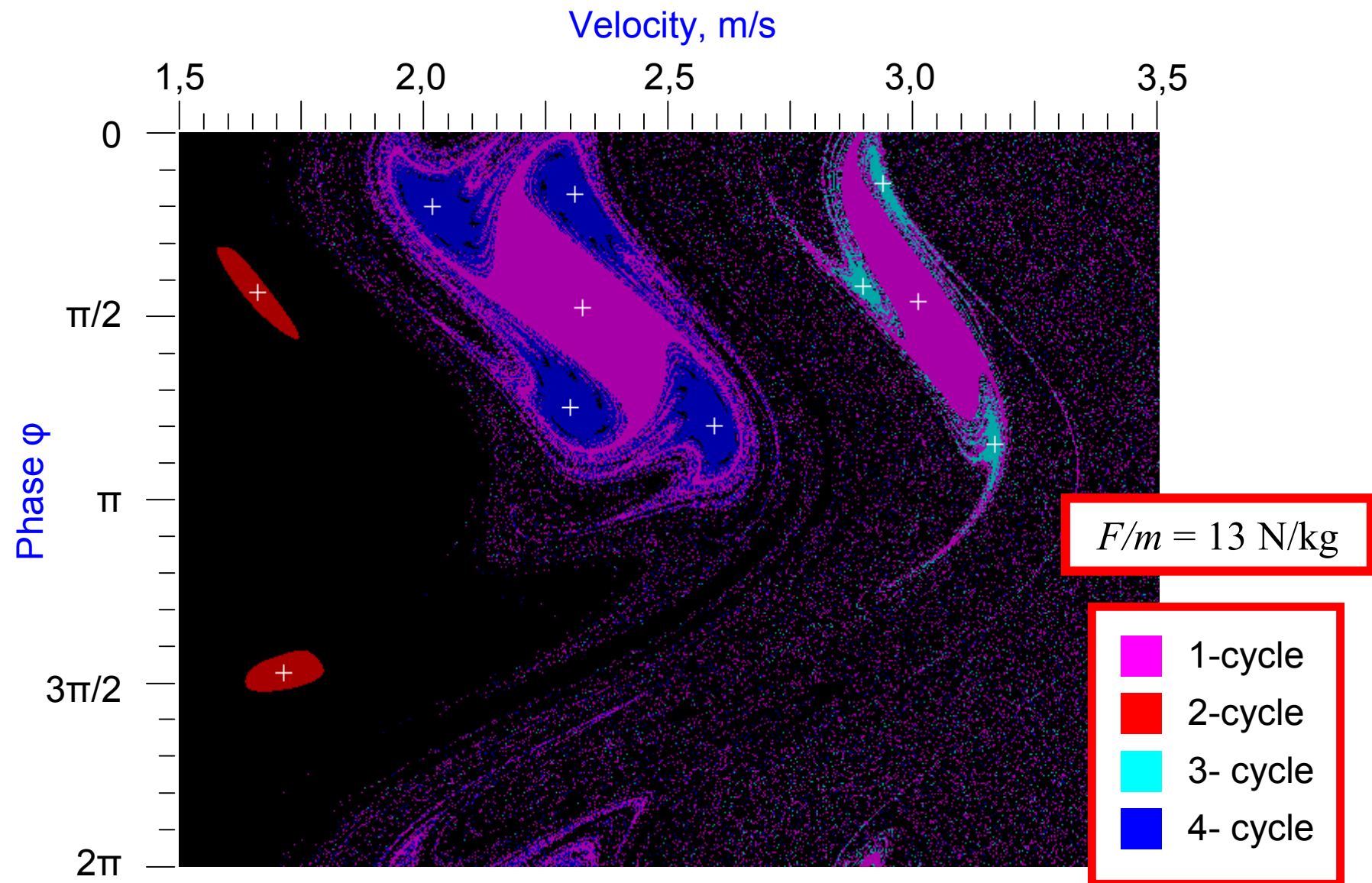
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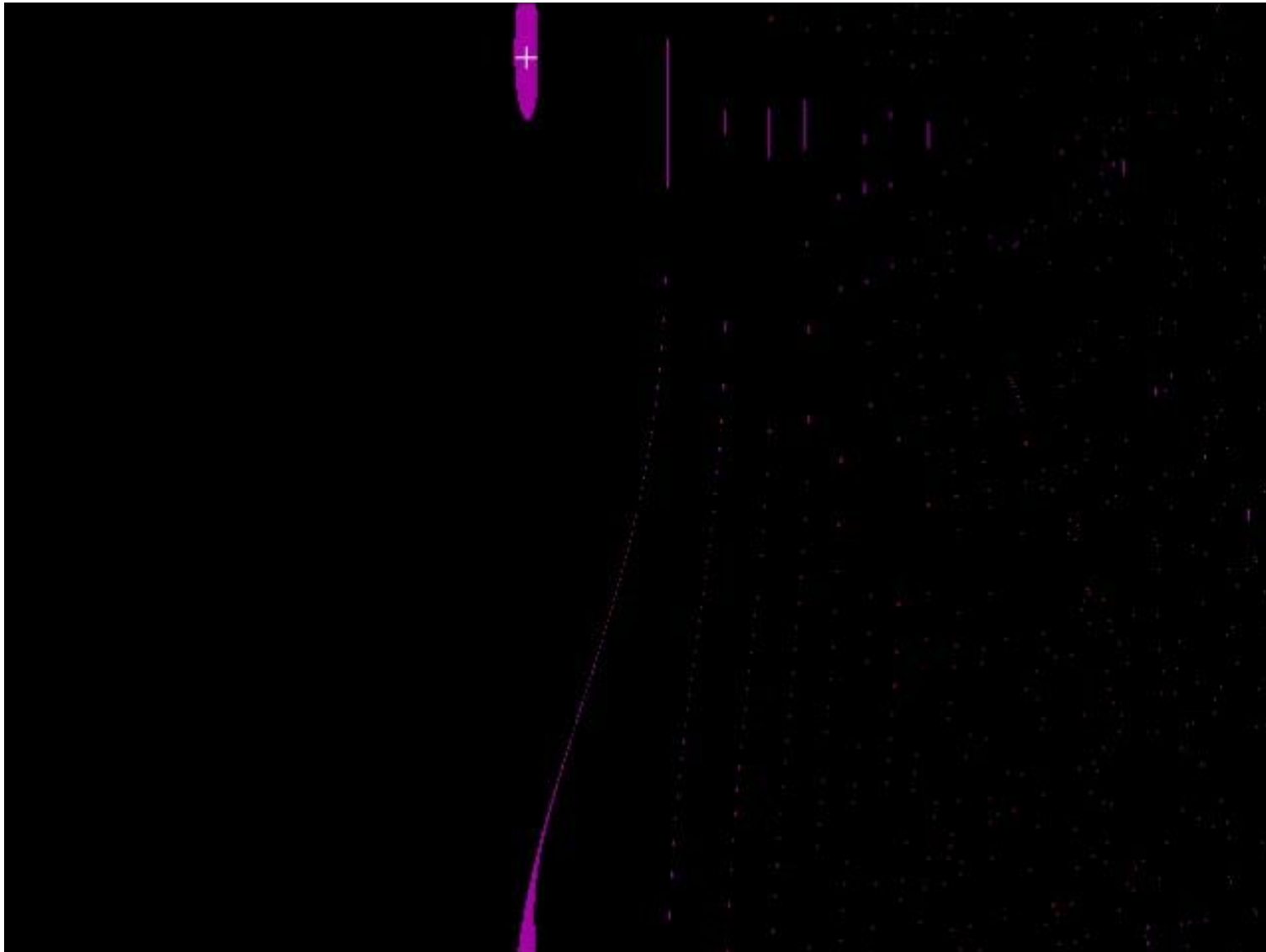
Poincare map

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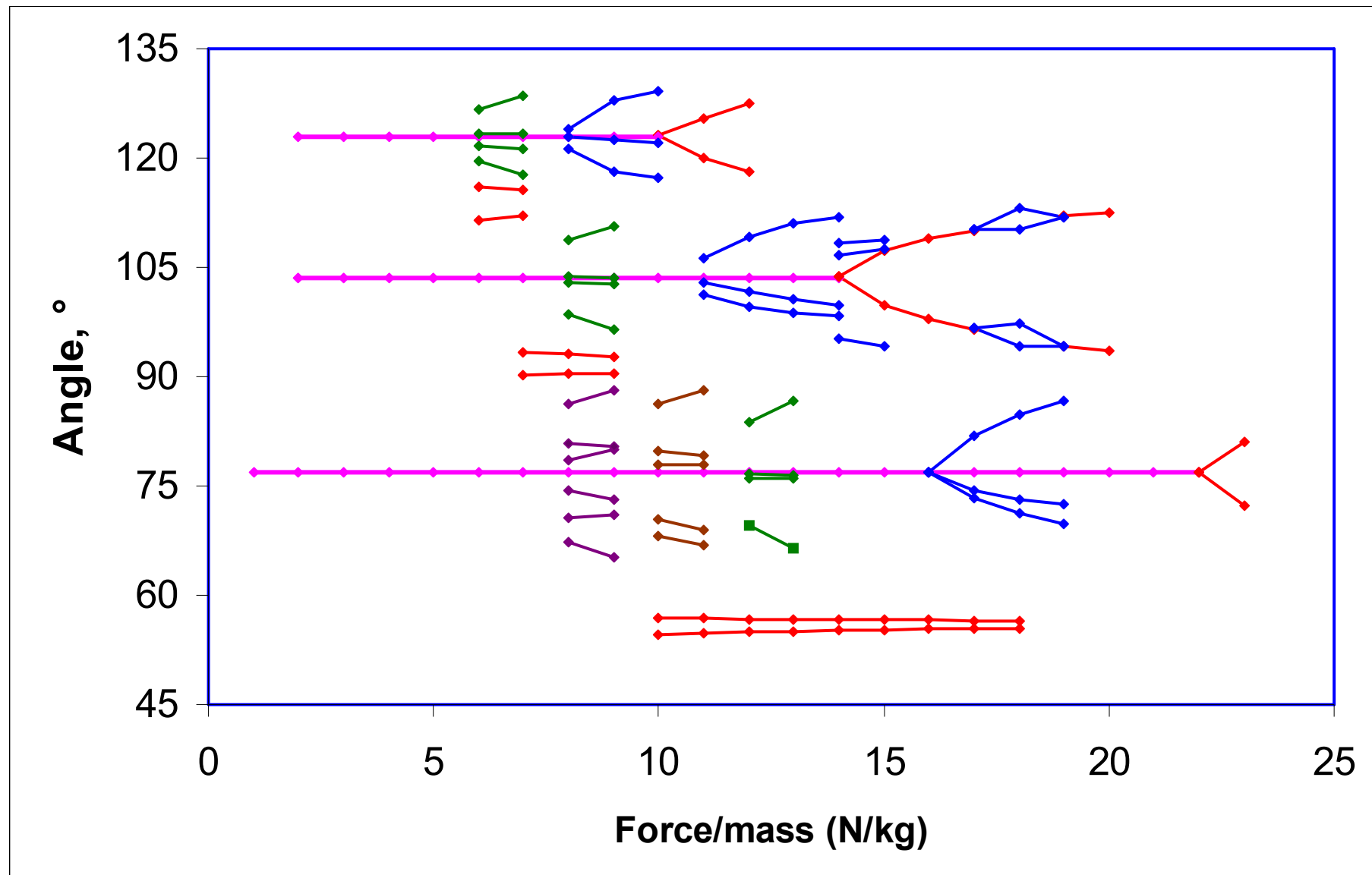
Increasing the strength (video)

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Behavior of the fixed points

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Summary

Conclusions

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AC solenoid

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Undamped oscillations (2 periods)

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Undamped oscillations

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3-periodic cycle

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Change in velocity for half-period

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Poincare mapping

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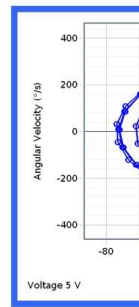
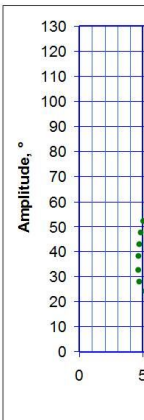
Poincare map

44

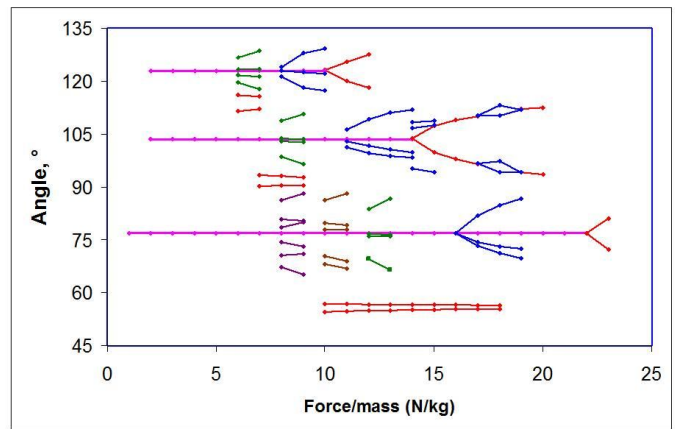
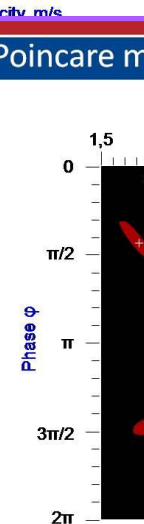
Behavior of the fixed points

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is braked



$$v_{n+1} = (1$$



- П.С.Ланда, Я.Б.Дубошинский (1989) “Автоколебательные системы с высокочастотными источниками энергии”. *УФН*, **158**, 729–742.
- V.Damgov, I.Popov (2000) “Discrete oscillations and multiple attractors in kick-excited systems”. *Discrete Dynamics in Nature and Society*, **4**, 99–124.



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**Thank you for
your attention!**