

Russia IYPT

Hovercraft

Nikolay Sibiryakov

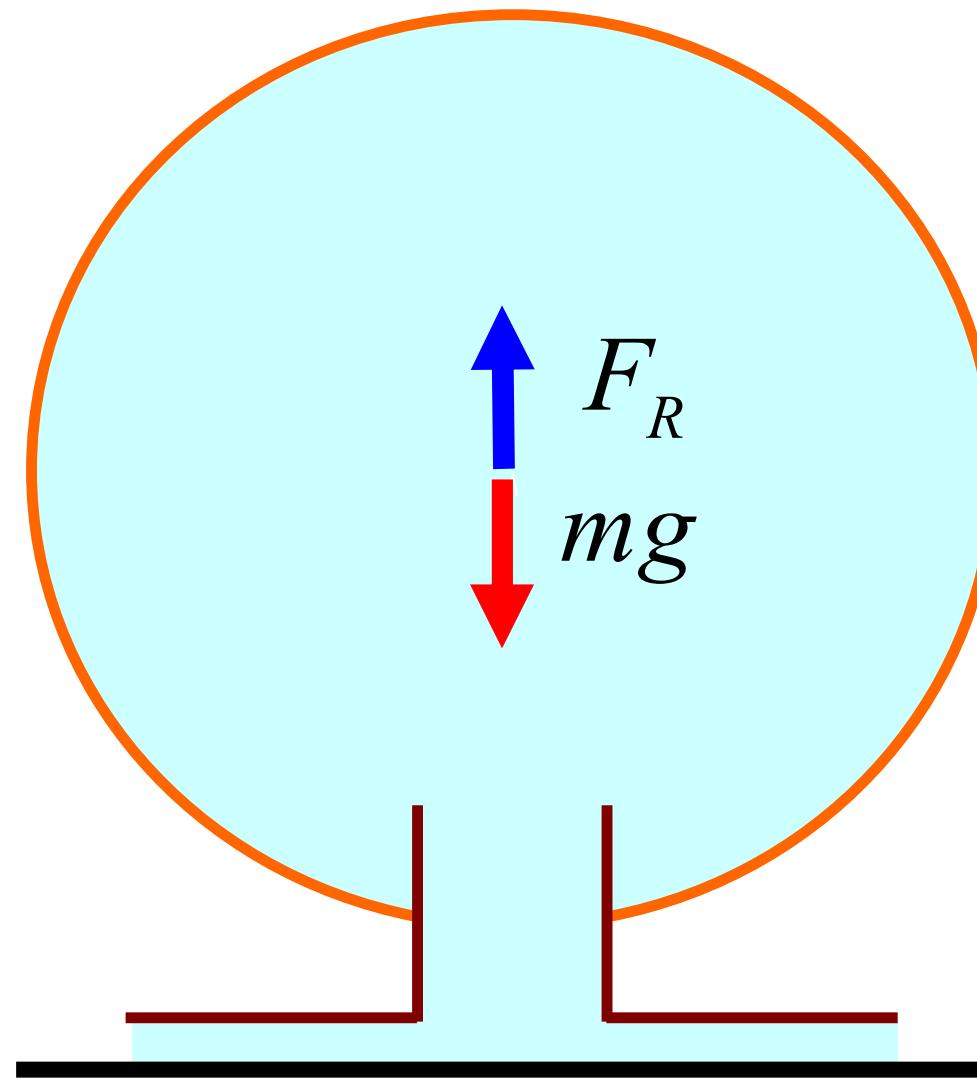
A simple model hovercraft can be built using a CD and a balloon filled with air attached via a tube. Exiting air can lift the device making it float over a surface with low friction. Investigate how the relevant parameters influence the time of the 'low-friction' state.

First observations

Floating over the table

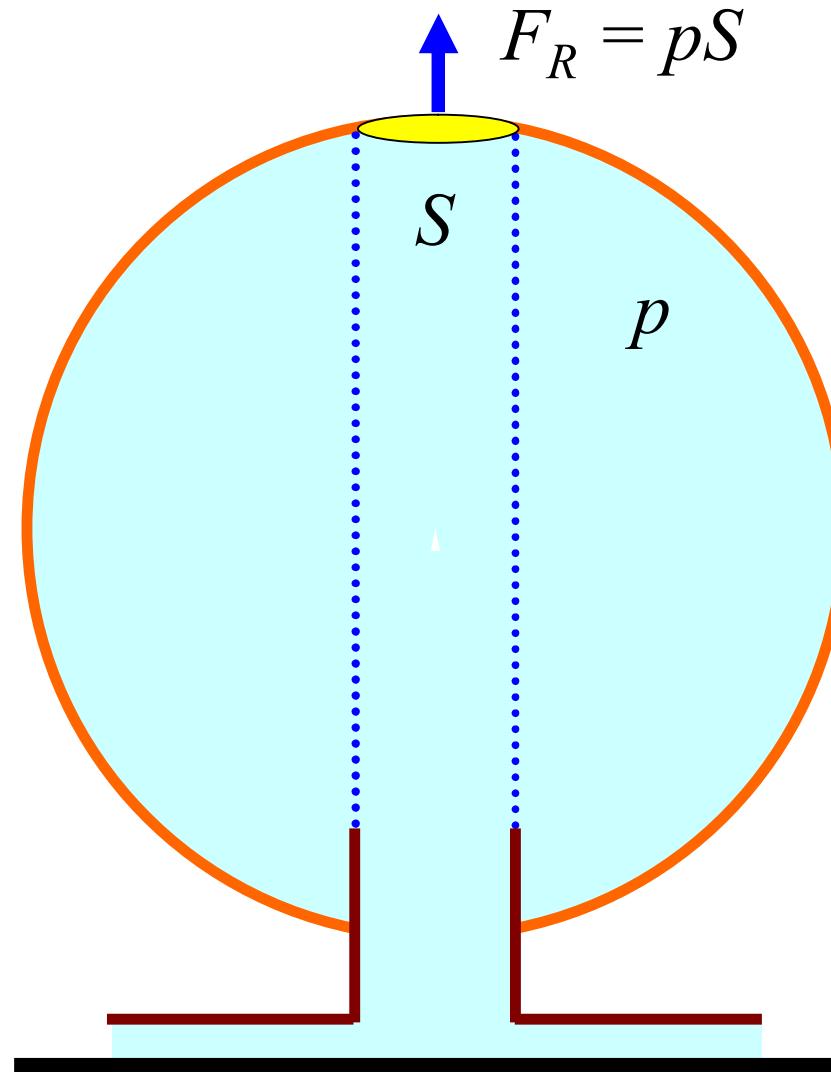
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The reactive force

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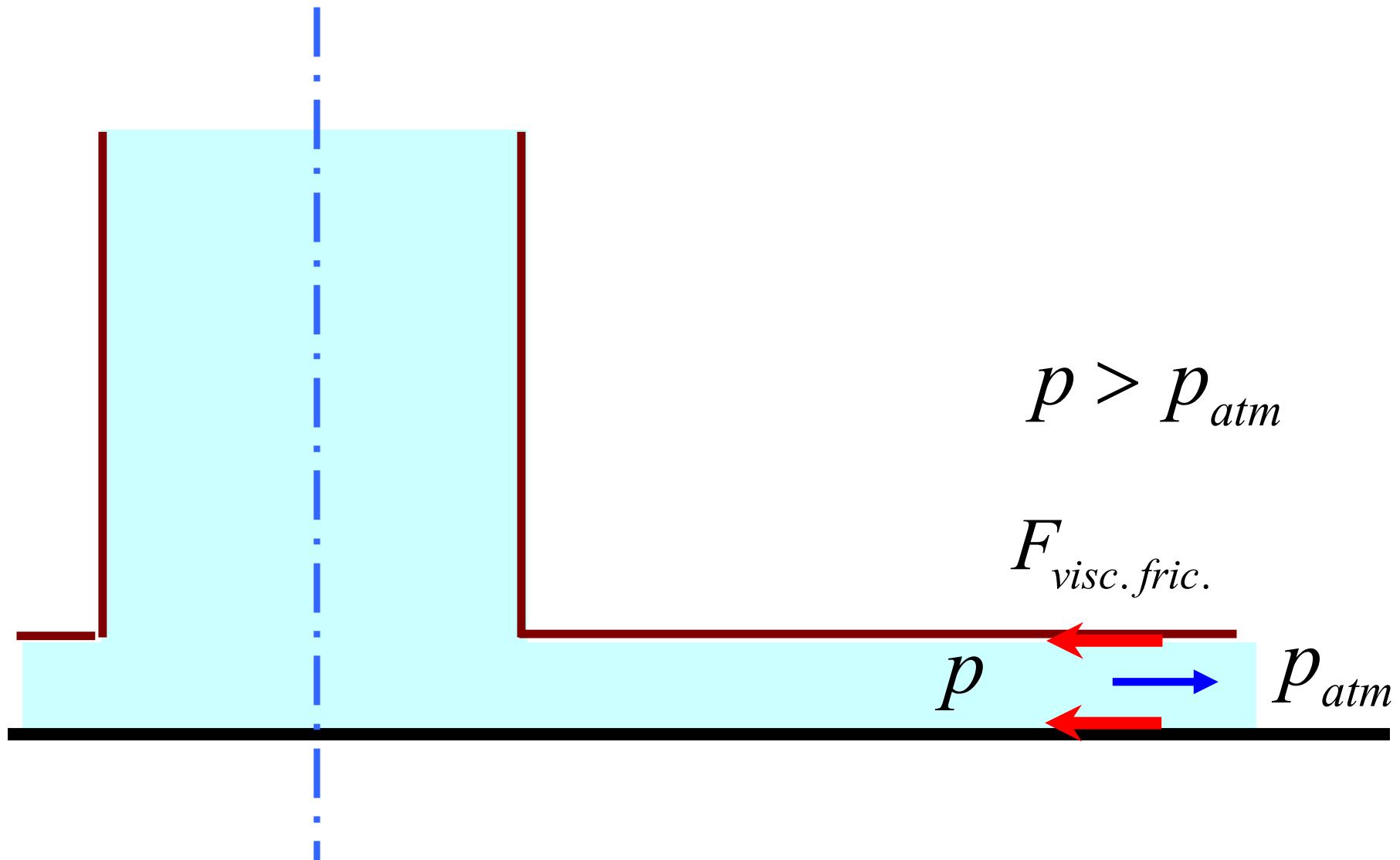


$$d = 13,8 \text{ mm}$$

$$S = 1,5 \text{ cm}^2$$

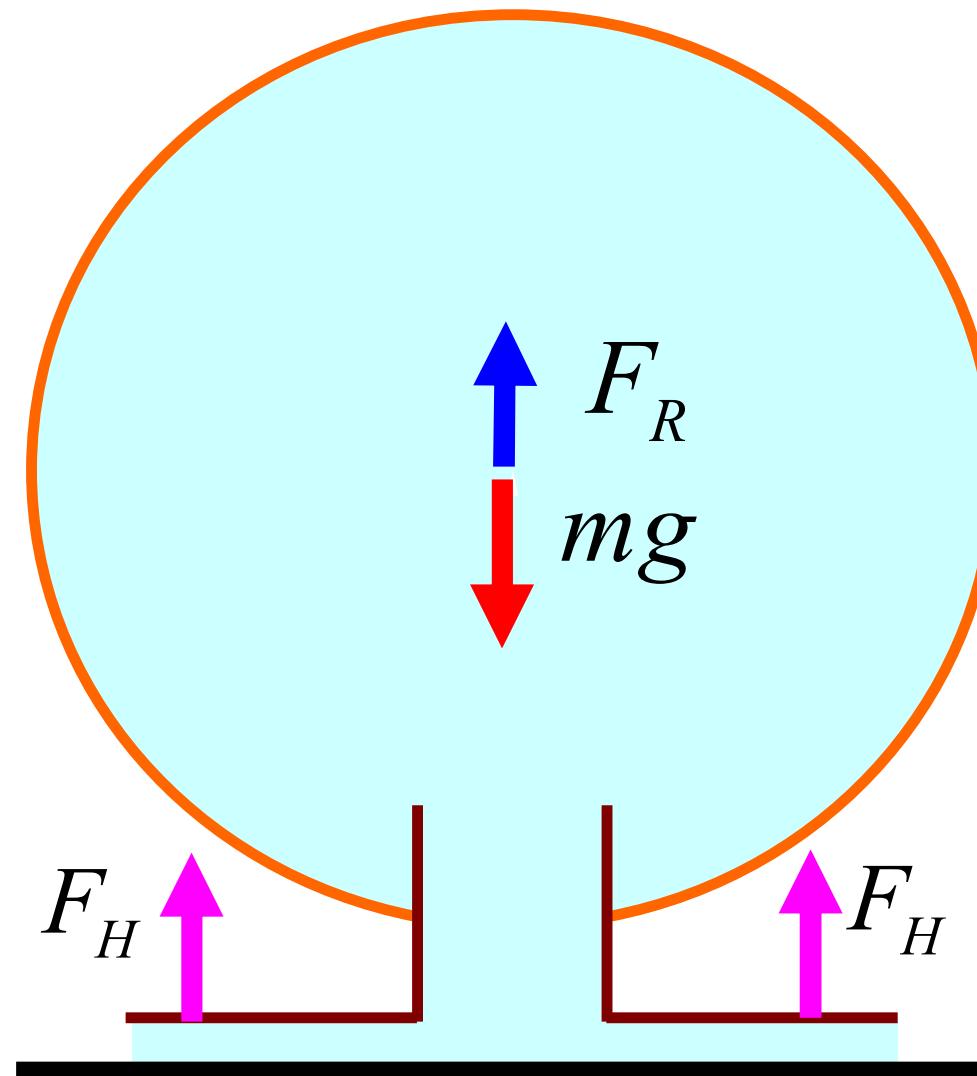
$$F_R = 0,15 \text{ N}$$

$$mg = 0,28 \text{ N}$$



Force balance over the table

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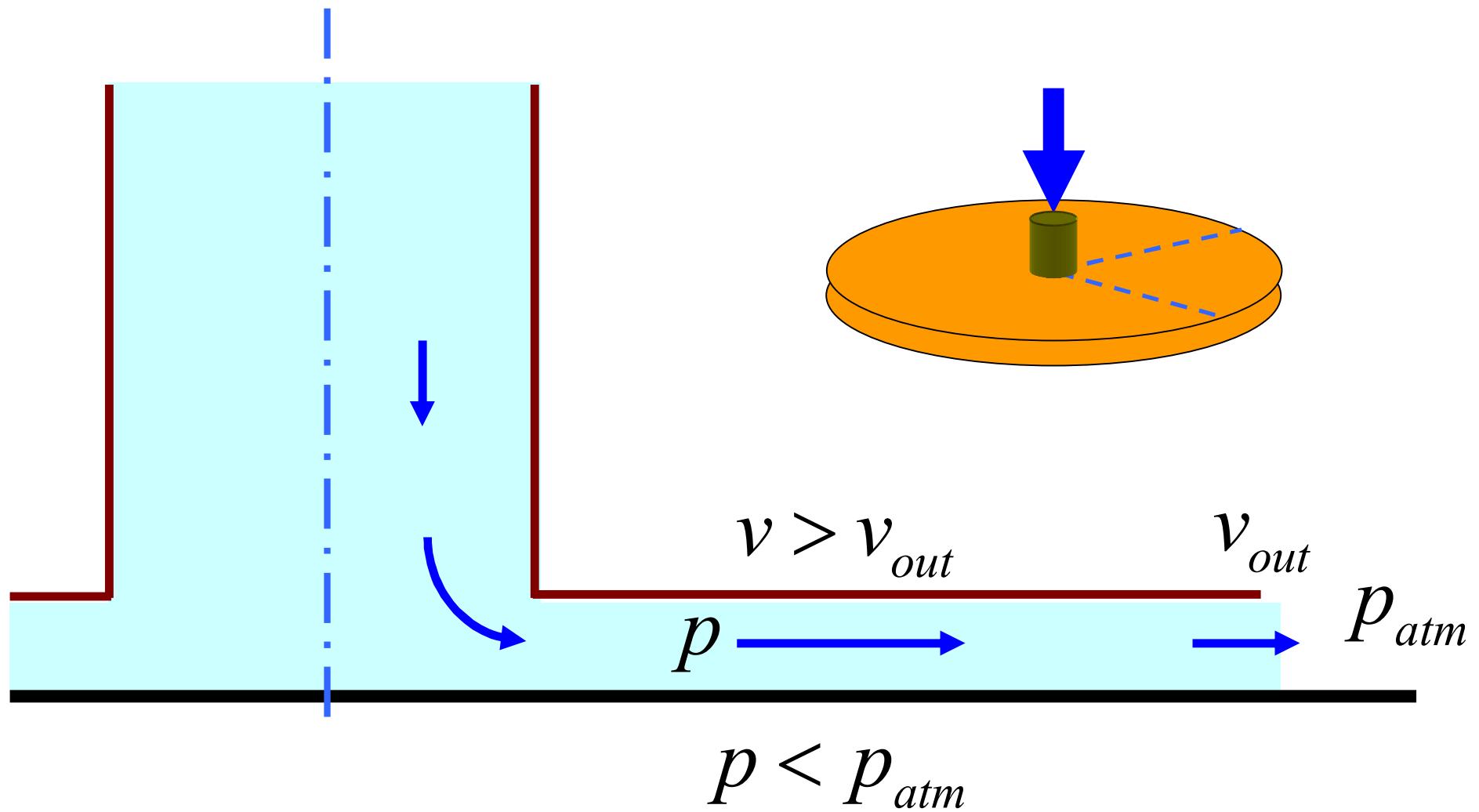
Floating under the ceiling

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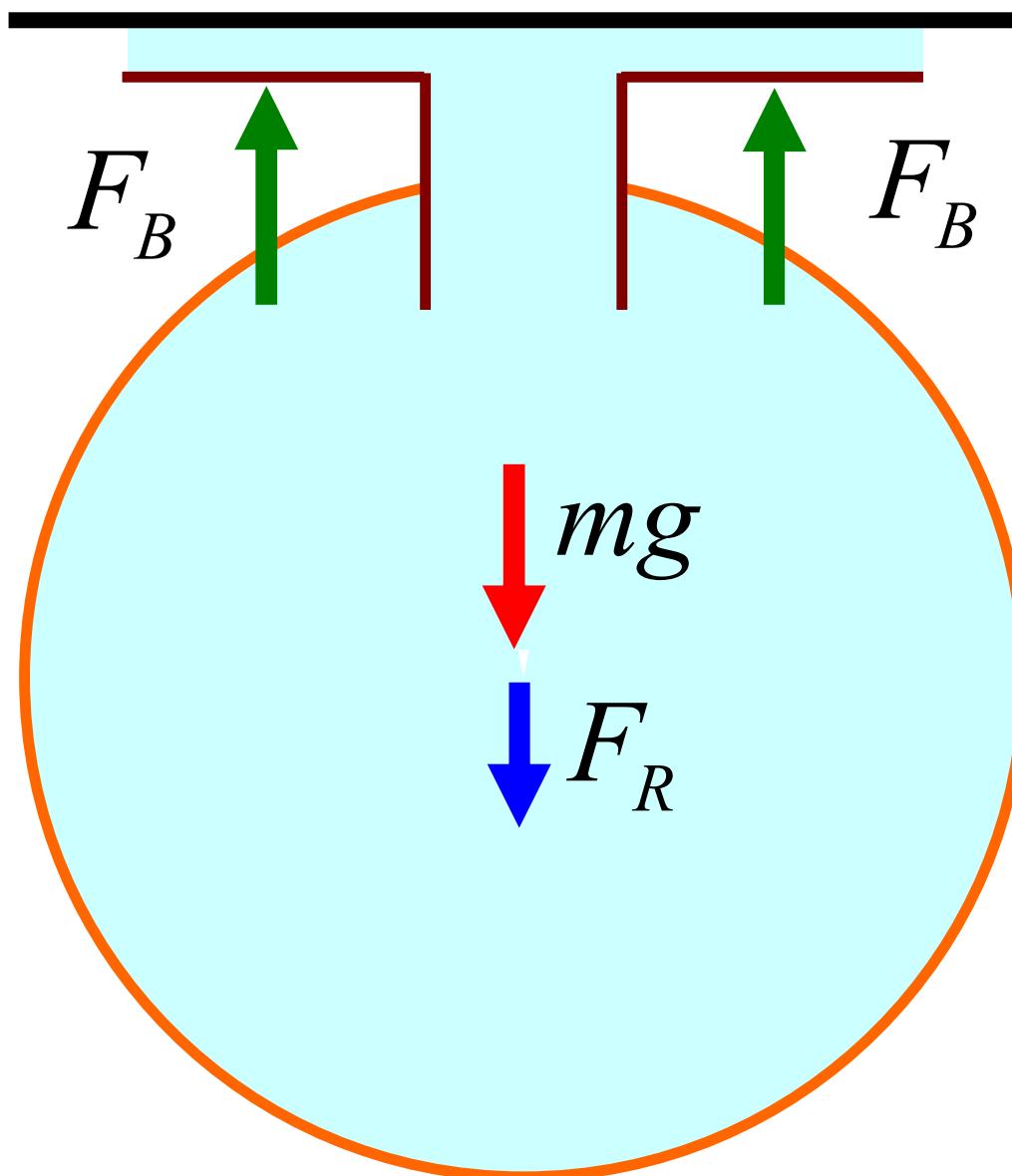
Bernoulli's principle

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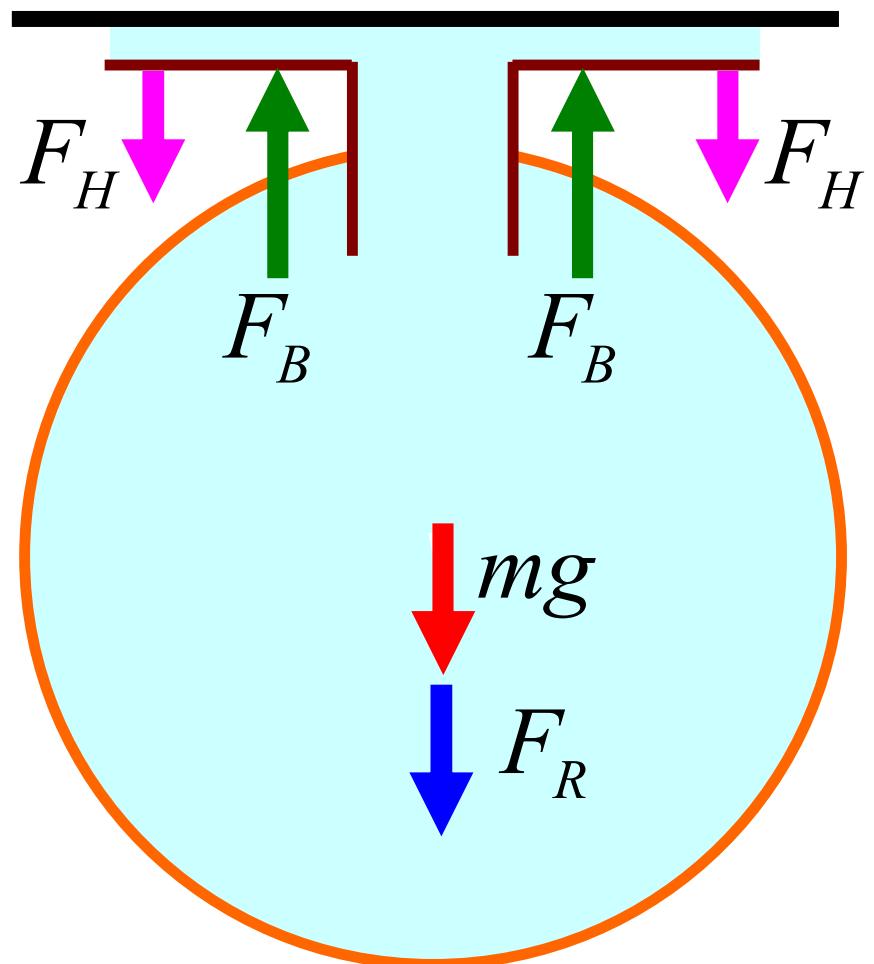
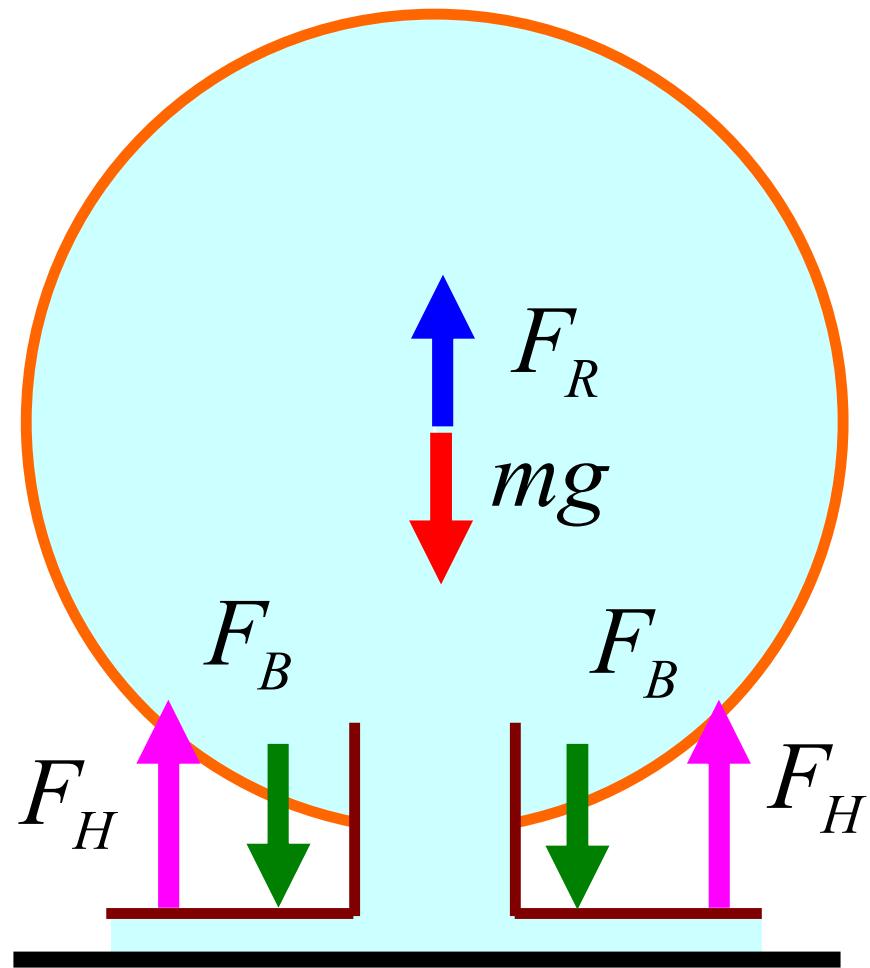
Force balance under the ceiling

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Final force balance

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Pressure under the disk



Hovering experiments



Theoretical model



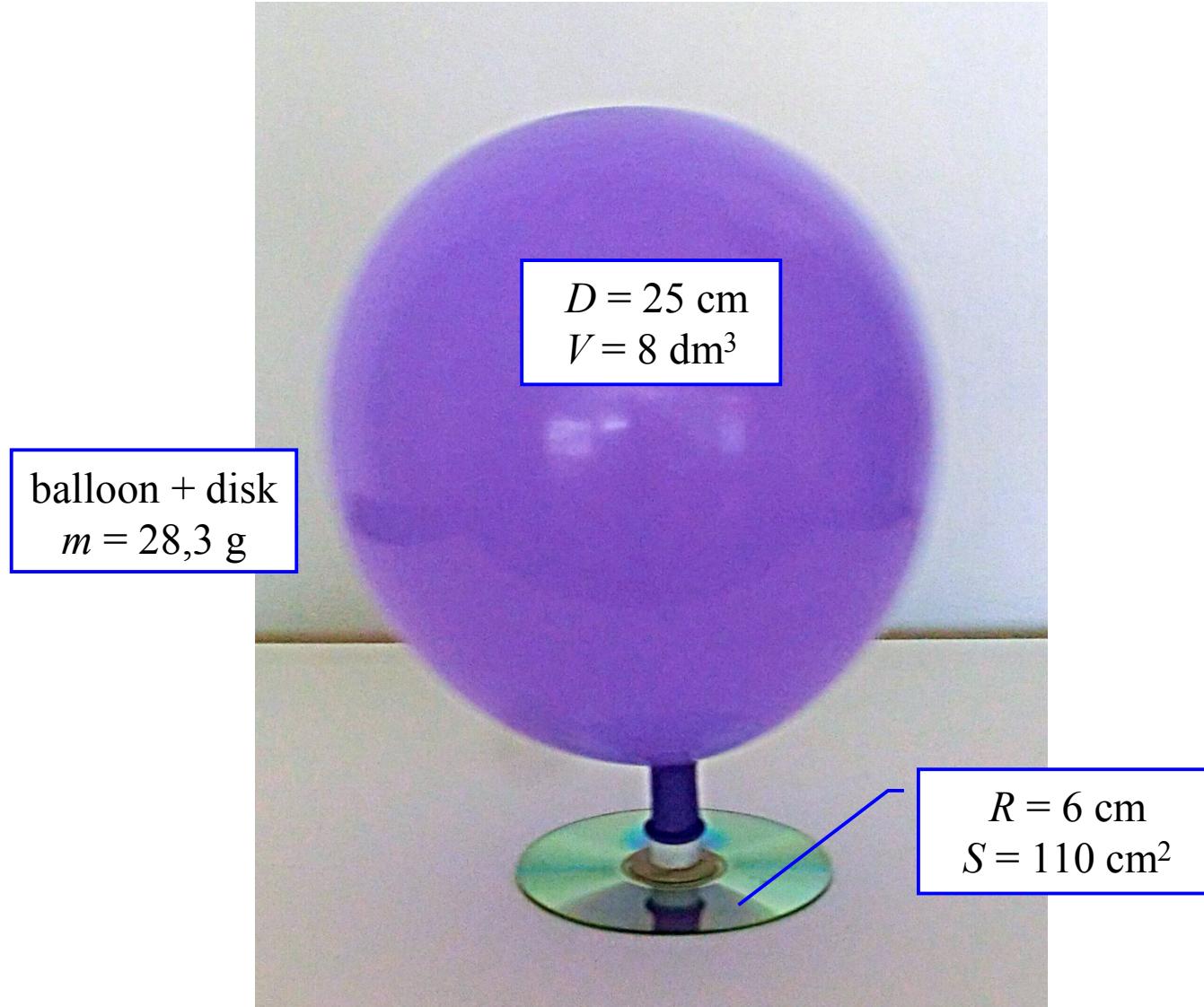
Hovering time vs. craft's weight

...how the relevant parameters influence the time of the 'low-friction' state...

Air pressure
under the disk

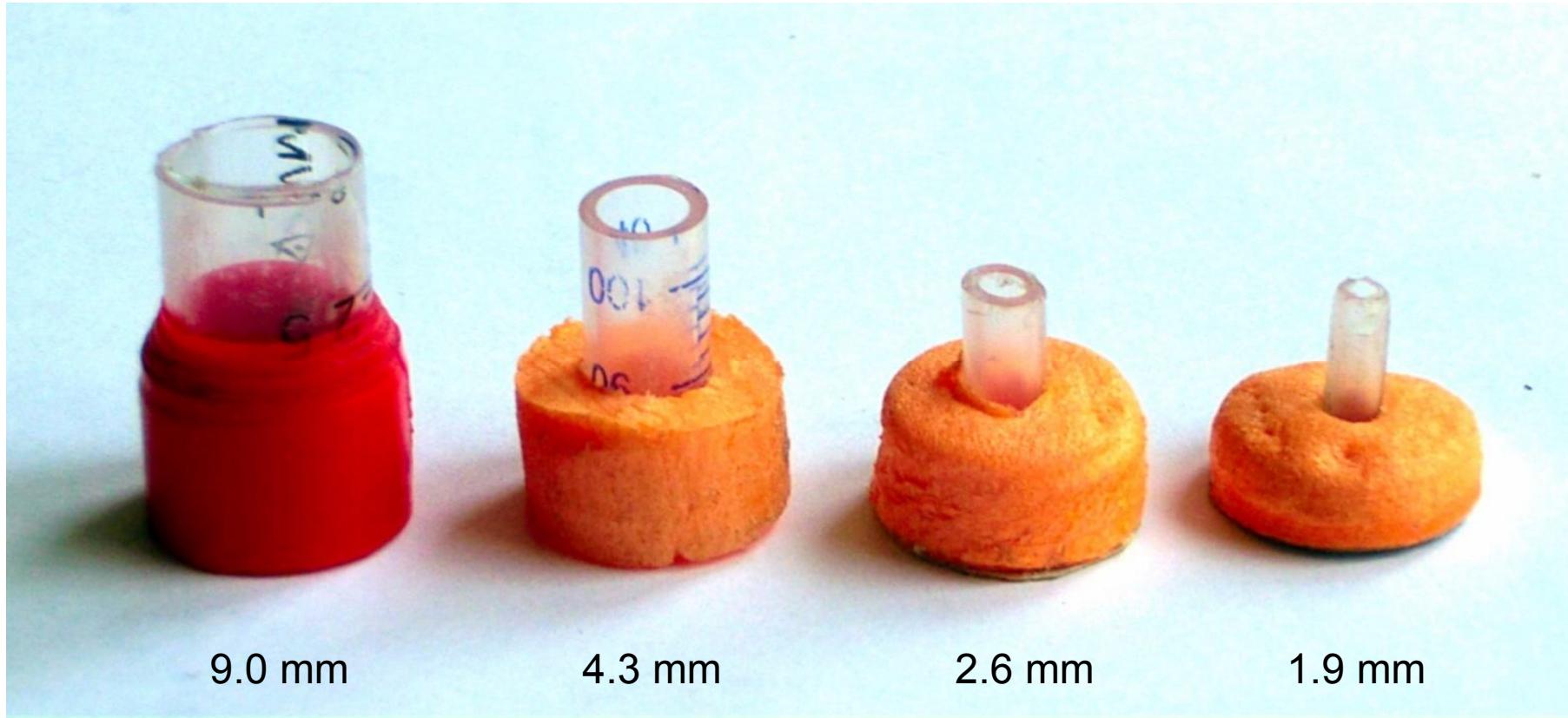
Parameters of the hovercraft

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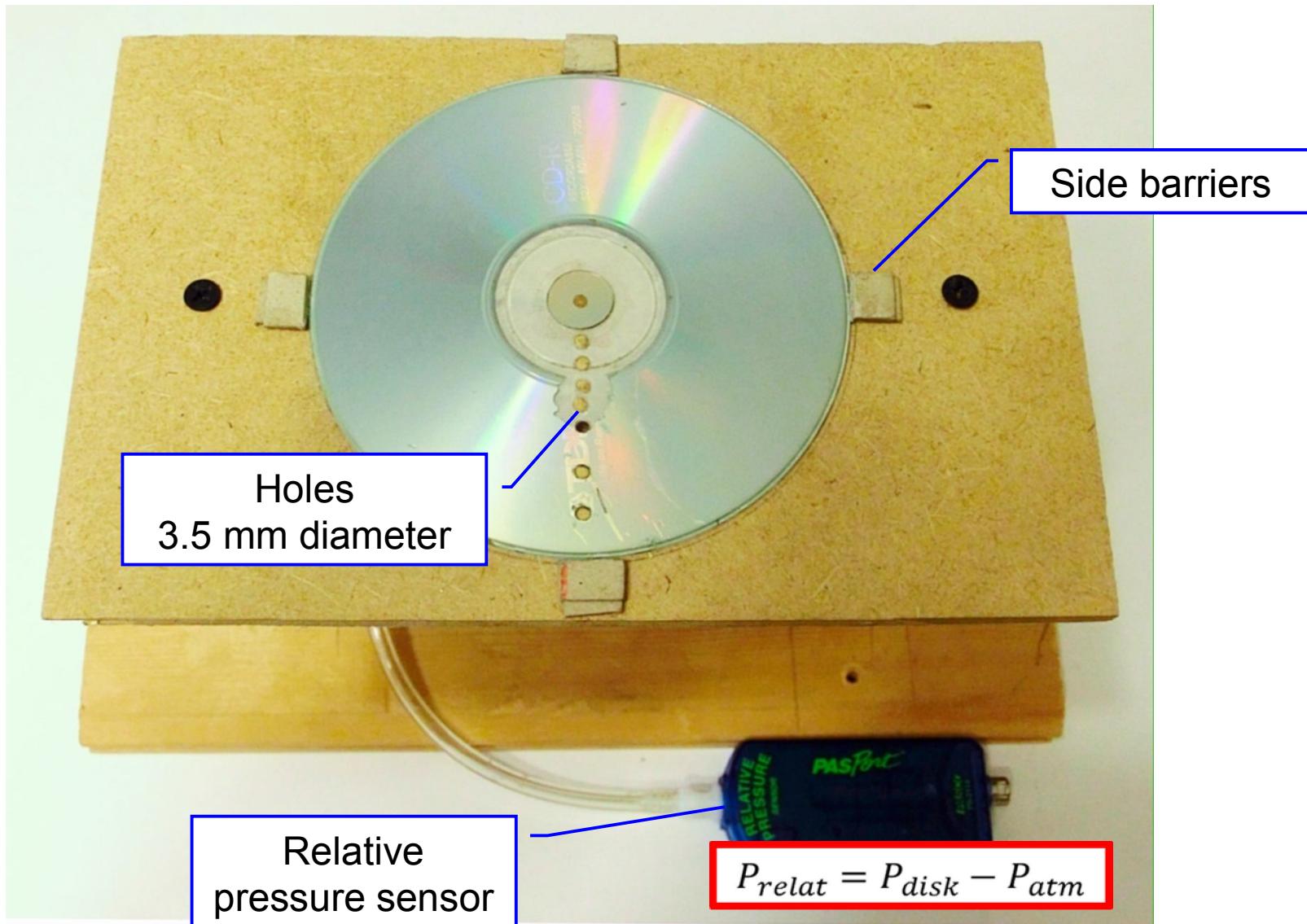
Plug-in nozzles

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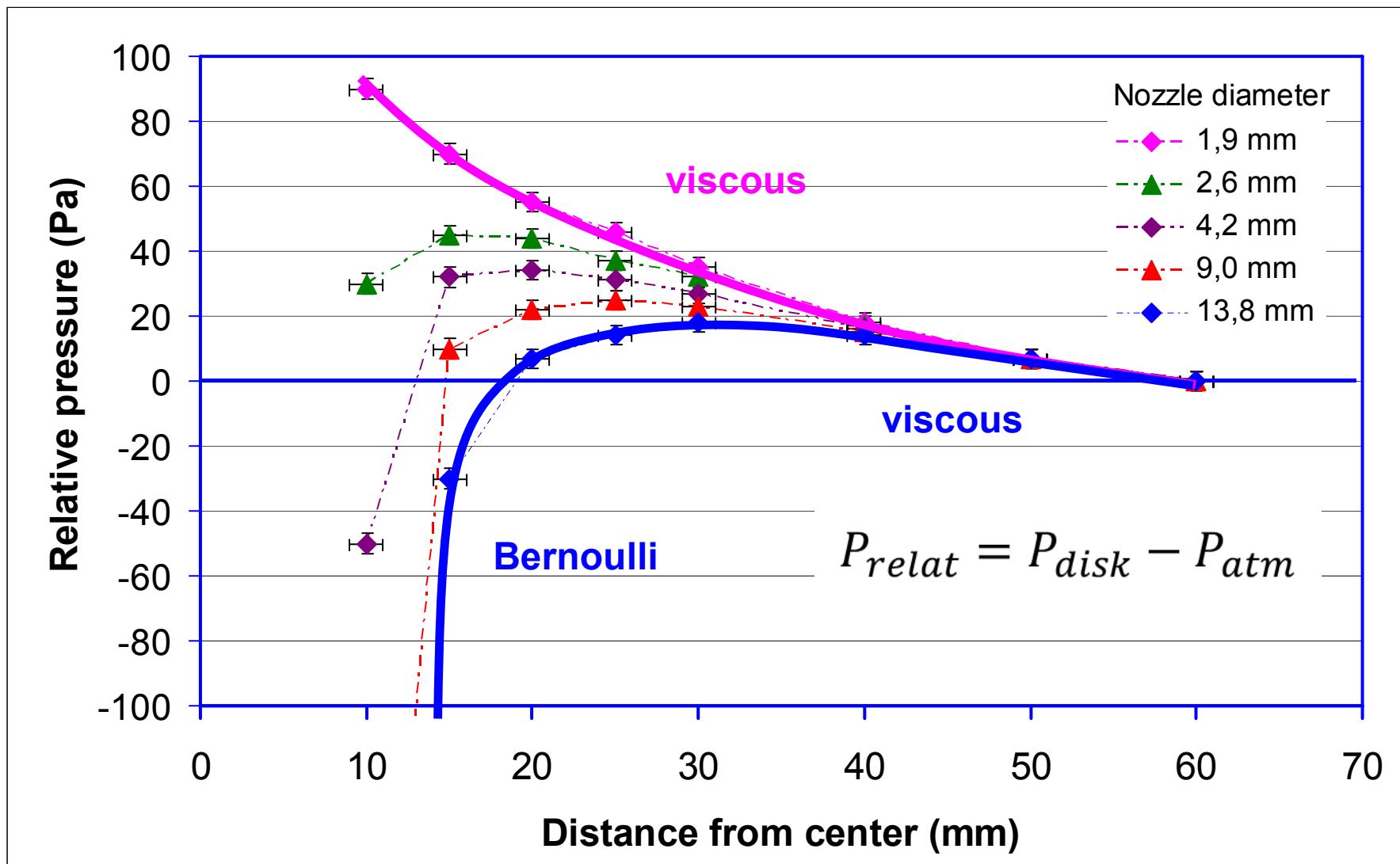
Setup for pressure measurement

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Pressure distribution under the disk

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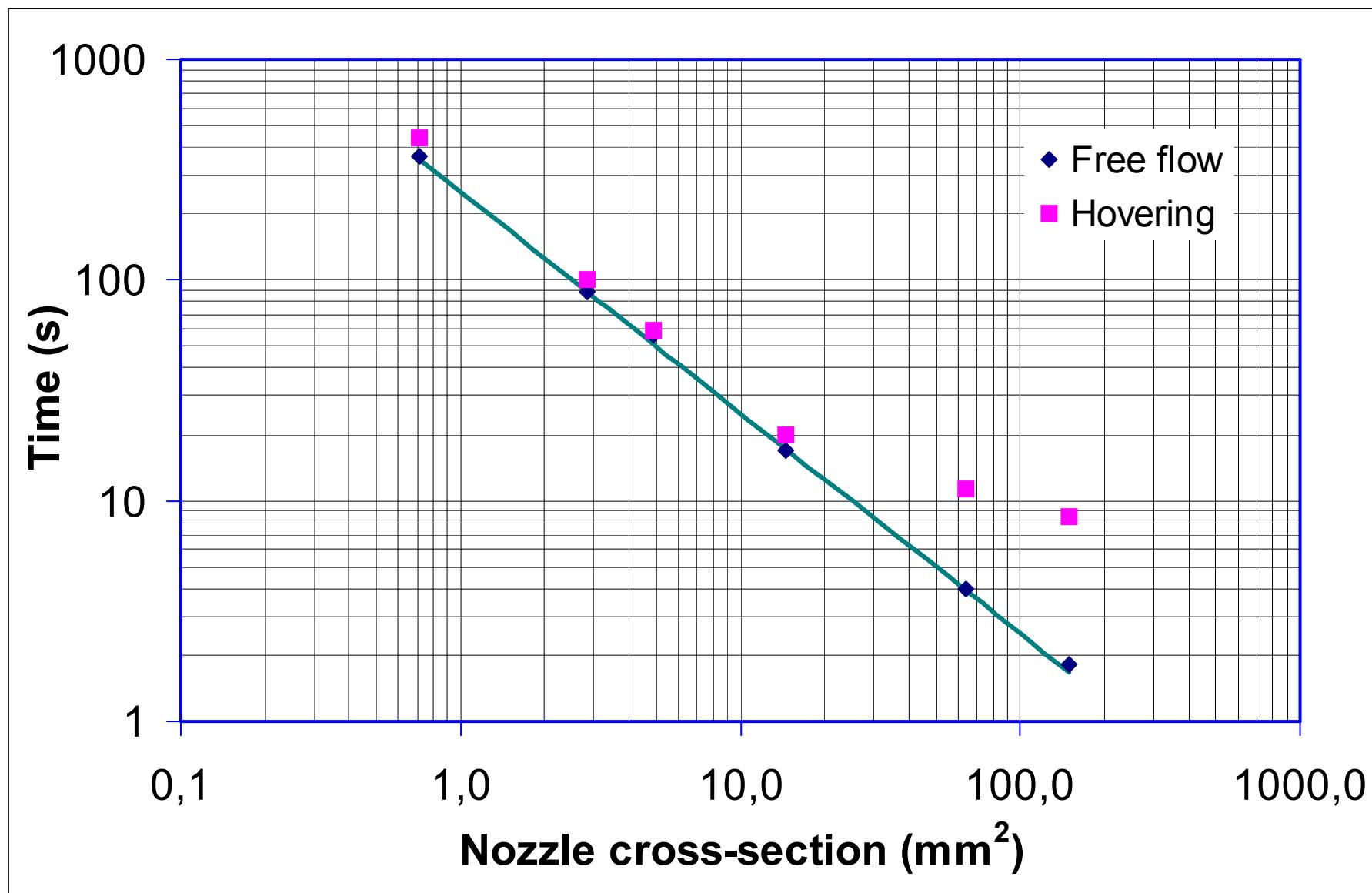


Hovering experiments

“...how the relevant parameters influence
the time of the 'low-friction' state...”

Time vs. nozzle cross-section

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Narrow nozzles
Viscosity dominates

What do we already know?

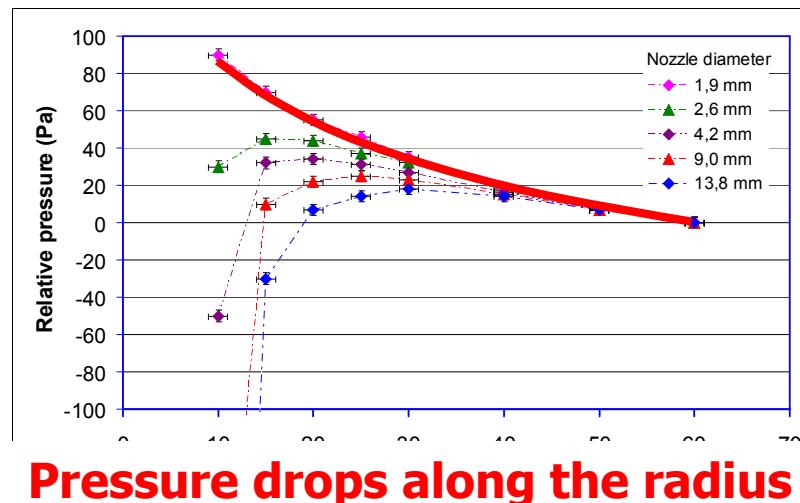
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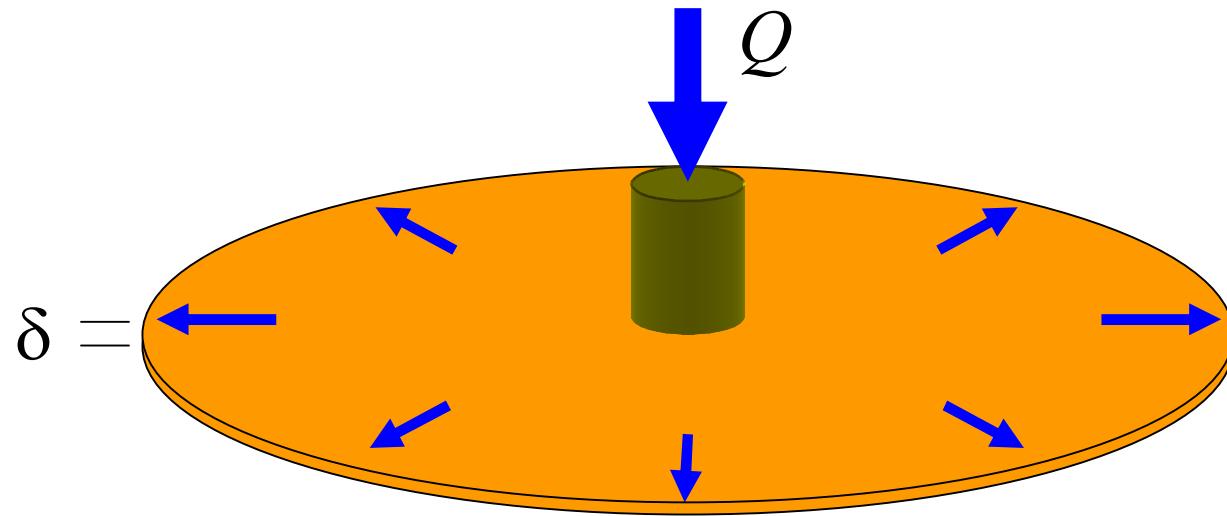
Small nozzle diameter



Almost free outflow



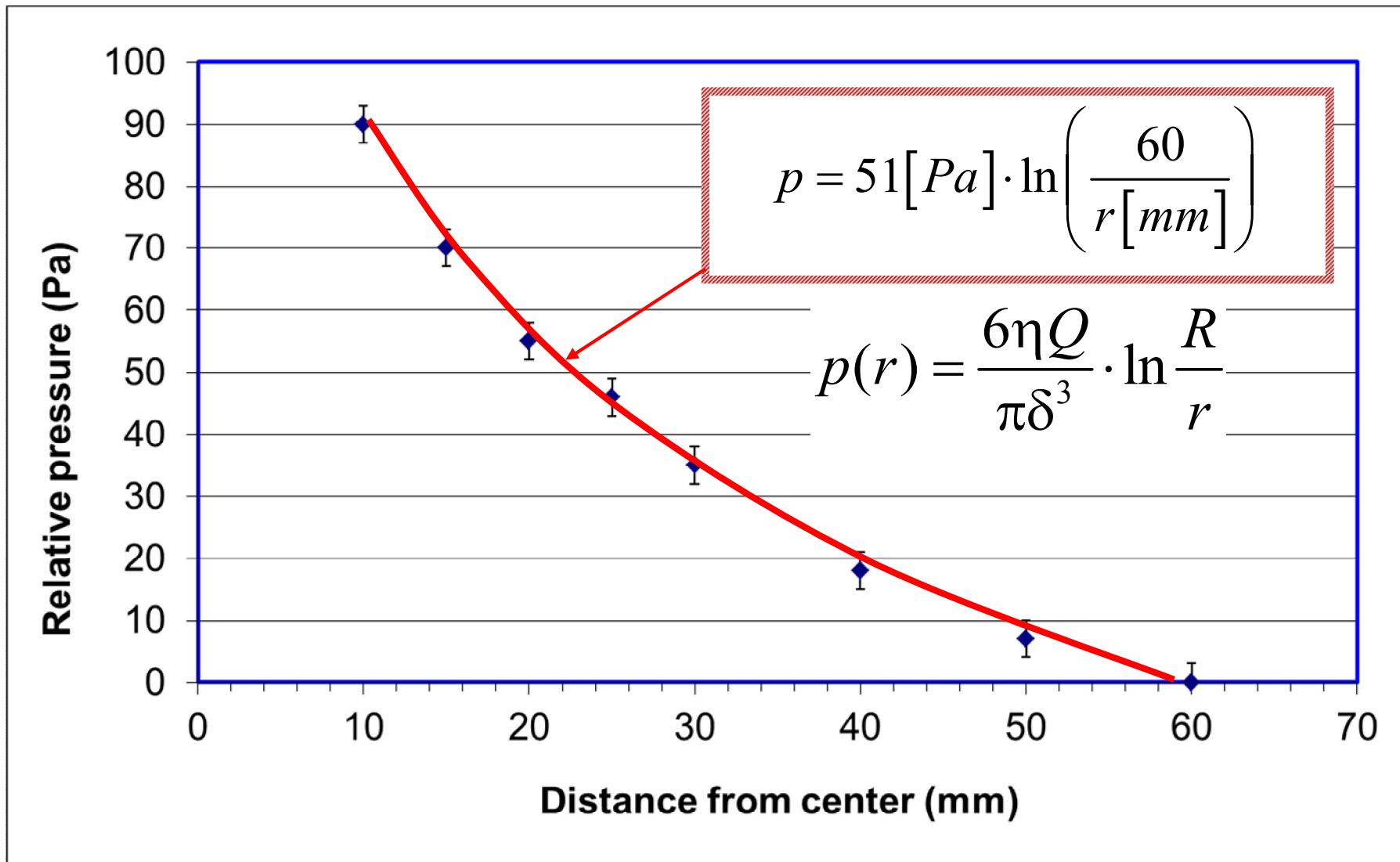
Pressure drops along the radius



Continuity condition: $v(r) = \frac{Q}{2\pi r \delta}$

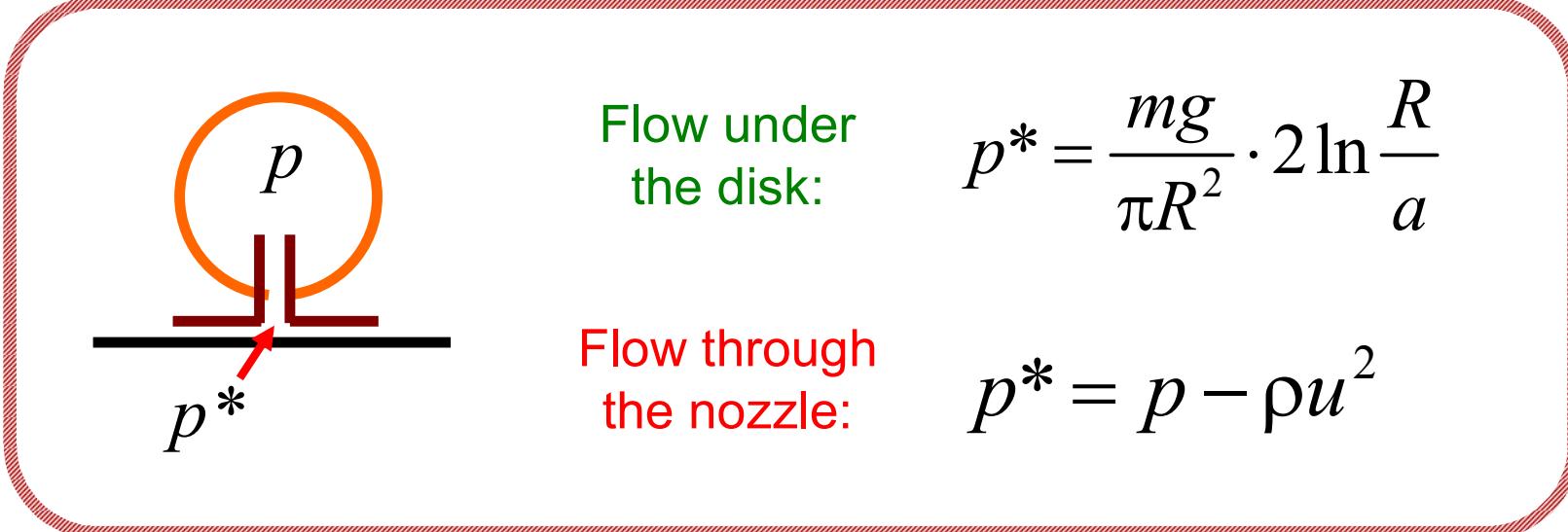
Darcy's law: $v(r) = -\frac{\delta^2}{12\eta} \frac{dp}{dr}$

$$\left. \begin{array}{l} v(r) = \frac{Q}{2\pi r \delta} \\ v(r) = -\frac{\delta^2}{12\eta} \frac{dp}{dr} \end{array} \right\} p(r) = \frac{6\eta Q}{\pi \delta^3} \cdot \ln \frac{R}{r}$$



Hovering time with narrow nozzle

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Velocity in the nozzle:

$$u^2 = \frac{p}{\rho} \cdot \left\{ 1 - \frac{mg}{p \cdot \pi R^2} \cdot 2 \ln \left(\frac{R}{a} \right) \right\}$$

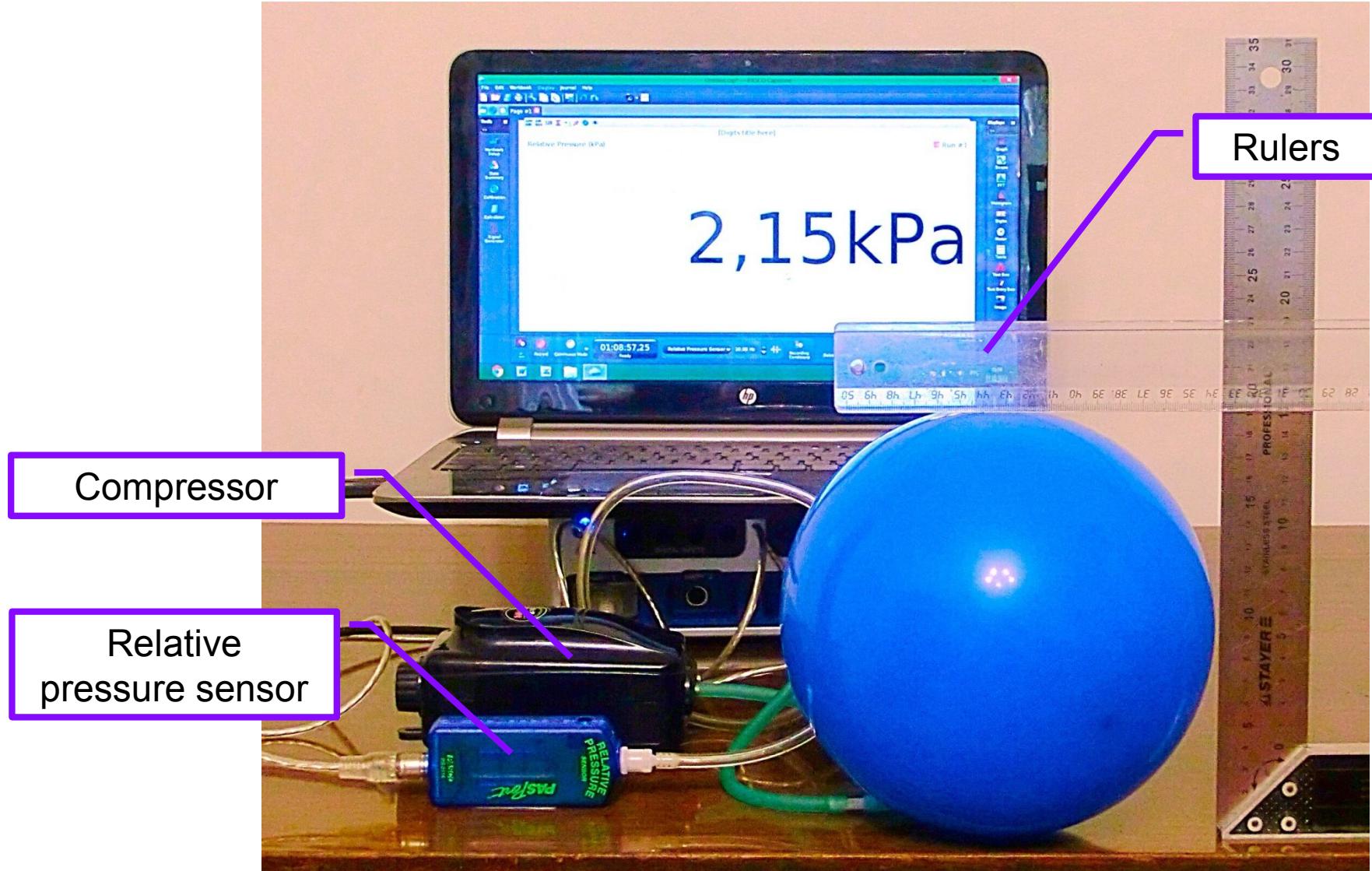
$$\tau = \frac{V}{S} \sqrt{\frac{\rho}{p}} \cdot \left(1 - \frac{mg}{p \cdot \pi R^2} \cdot 2 \ln \frac{R}{a} \right)^{-1/2}$$

The term $\frac{V}{S} \sqrt{\frac{\rho}{p}}$ is highlighted with a red box and a red arrow points to it from the top left. The term p in the denominator of the fraction $\frac{mg}{p \cdot \pi R^2}$ is also highlighted with a red box and a red arrow points to it from the bottom right.

We should know it!

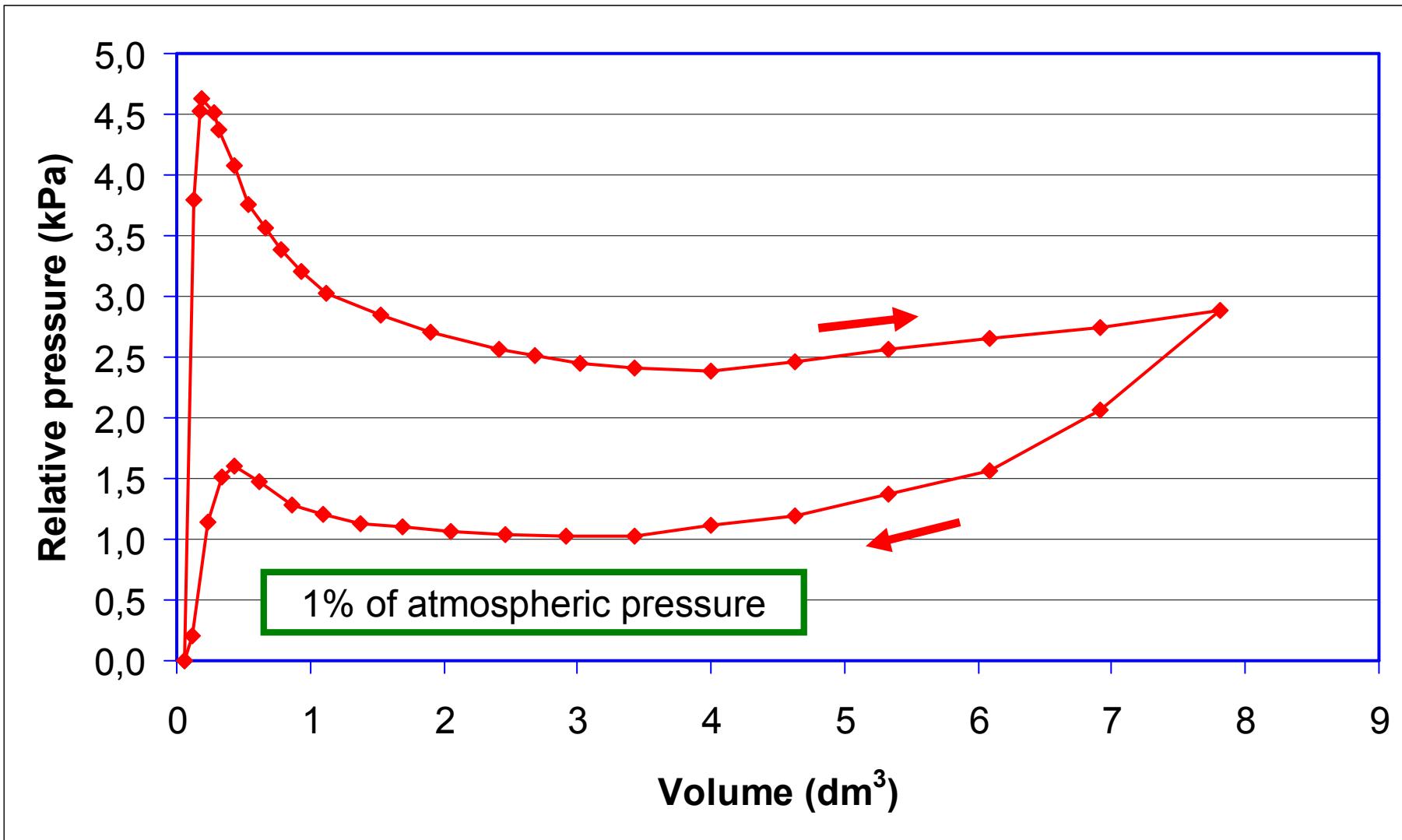
Measurement of relative pressure

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Pressure vs. volume

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$$\tau = \tau_0 \cdot \left(1 - \frac{mg}{p \cdot \pi R^2} \cdot 2 \ln \frac{R}{a} \right)^{-1/2}$$

For our parameters of the balloon and the disk

$$\tau \approx \tau_0 \left(1 + 0,02 \cdot \ln \frac{R}{a} \right) \quad +$$

If $R/a = 60$, $\tau = 1,08 \cdot \tau_0$

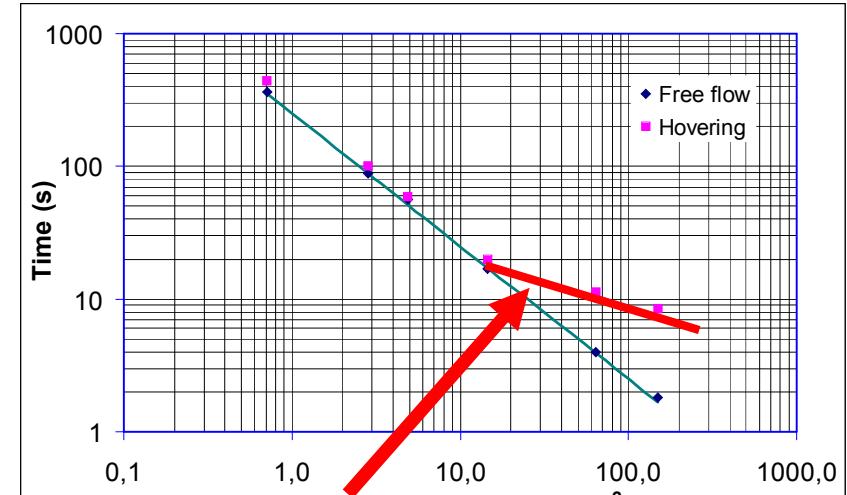
Wide nozzles
Viscosity + Bernoulli

What do we already know?

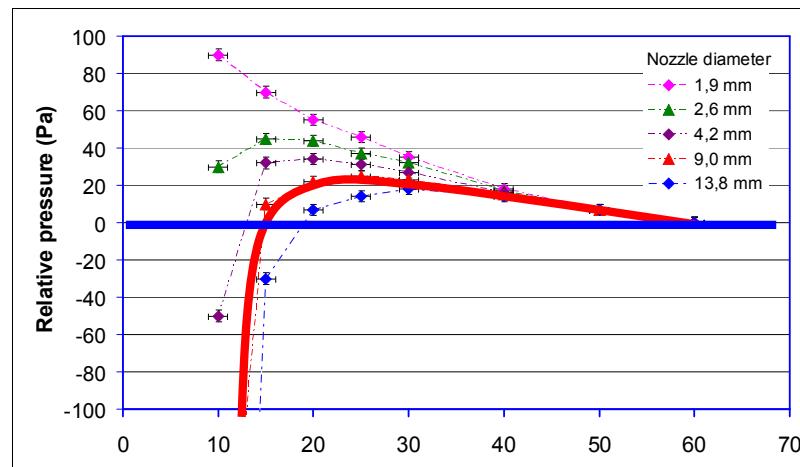
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Big nozzle diameters



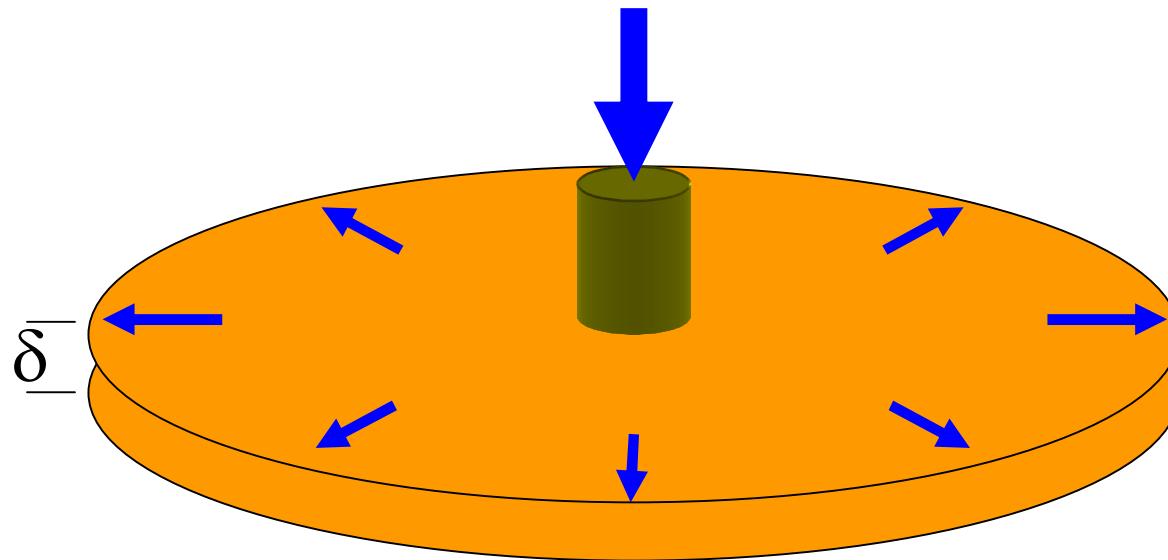
Outflow differs from the free outflow



Pressure under the central area is less than atmospheric

Wide gap: Bernoulli's regime

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Continuity condition: $v(r) = \frac{Q}{2\pi r \delta}$

Bernoulli's principle: $p + \frac{\rho v^2}{2} = \text{const}$

Relative Pressure is negative

$$p(r) = \frac{1}{4\pi^2} \cdot \frac{\rho Q^2}{\delta^2} \left(\frac{1}{R^2} - \frac{1}{r^2} \right)$$

$$p(r) = \underbrace{\frac{6\eta Q}{\pi \delta^3} \cdot \ln \frac{R}{r}}_{\text{Viscous}} + \underbrace{\frac{1}{4\pi^2} \cdot \frac{\rho Q^2}{\delta^2} \left(\frac{1}{R^2} - \frac{1}{r^2} \right)}_{\text{Bernoulli}}$$

$$p(r) = \underbrace{\frac{6\eta Q}{\pi \delta^3} \cdot \ln \frac{R}{r}}_{\text{viscous}} + \underbrace{\frac{27}{140\pi^2} \cdot \frac{\rho Q^2}{\delta^2} \left(\frac{1}{R^2} - \frac{1}{r^2} \right)}_{\text{Bernoulli with a parabolic profile}}$$

Armengol J., Calbó J., Pujol T., Roura P. (2011) “Bernoulli correction to viscous losses: Radial flow between two parallel discs”.

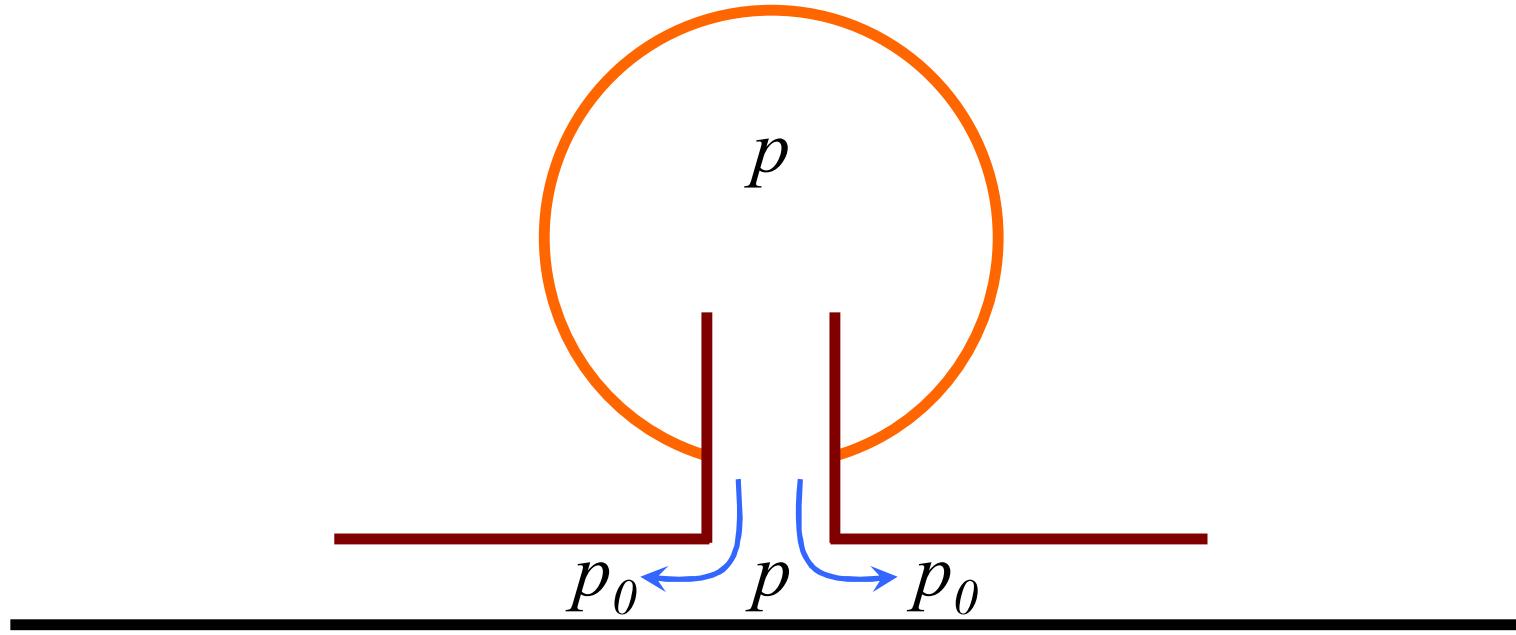
Force balance

$$mg - F_R \approx \underbrace{\frac{3\eta QR^2}{\delta^3}}_{\text{Viscous}} - \underbrace{\frac{27\rho Q^2}{70\pi\delta^2} \ln\left(\frac{R}{a}\right)}_{\text{Bernoulli}}$$

Both terms are large compared with the weight, so they are approximately equal to each other

$$Q \cdot \delta \approx \frac{70}{9} \cdot \frac{\eta \cdot \pi R^2}{\rho \ln(R/a)}$$

Outflow from the nozzle under the disk



As outflow from the nozzle into the atmosphere

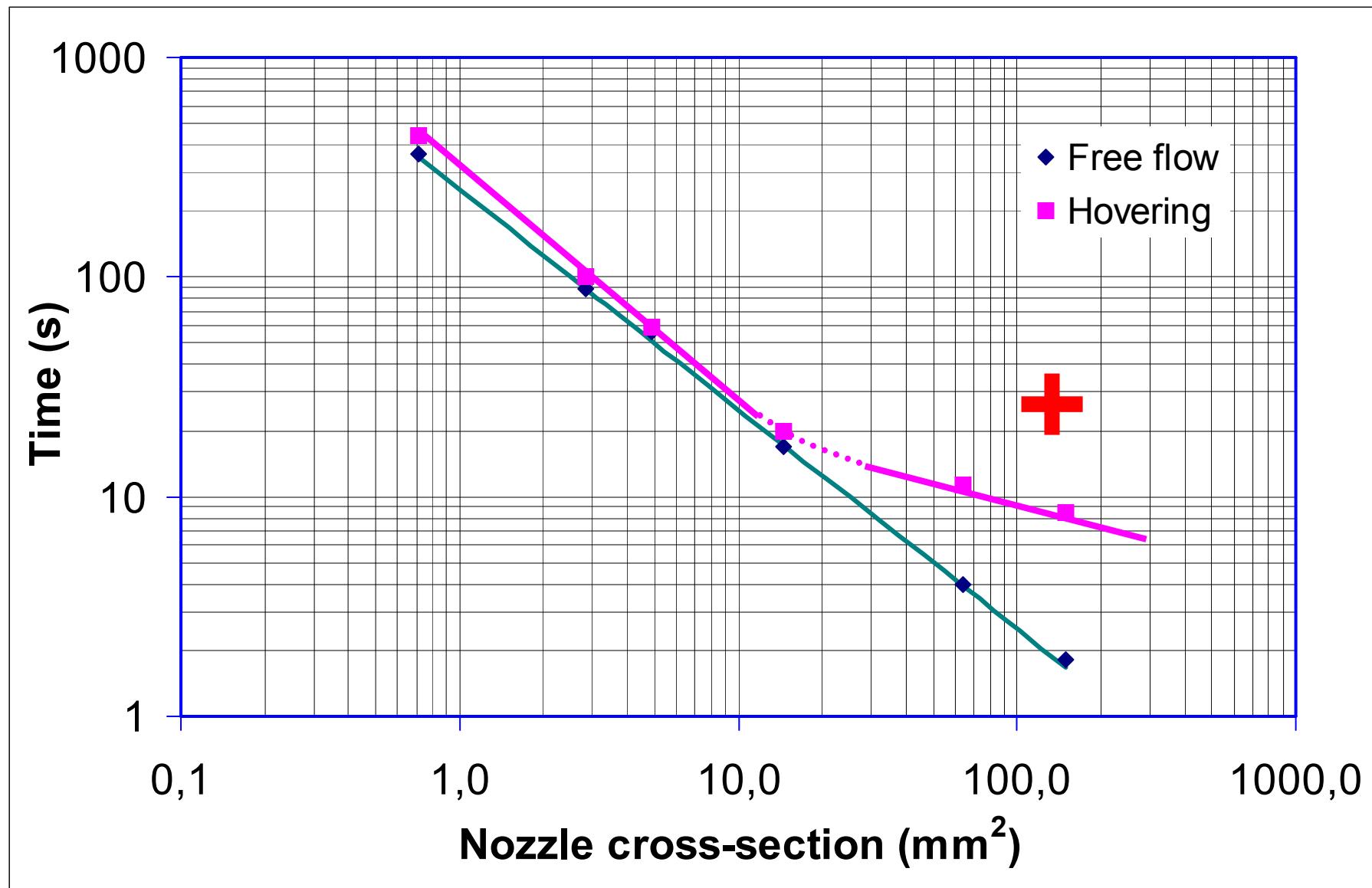
$$Q = \underbrace{2\pi a \delta}_{\text{Cylindrical entry}} \cdot v_0$$

$$v_0 = \sqrt{\frac{p}{\rho}}$$

Cylindrical entry

$$Q \cdot \delta \approx \frac{70}{9} \cdot \frac{\eta \cdot \pi R^2}{\rho \ln(R/a)} \quad Q = 2\pi a \delta \cdot v_0$$

$$\tau = V \cdot \sqrt{\frac{9}{70} \cdot \frac{\rho \ln(R/a)}{2\pi a \cdot \pi R^2 \cdot \eta v_0}}$$



Hovering time
vs.
weight

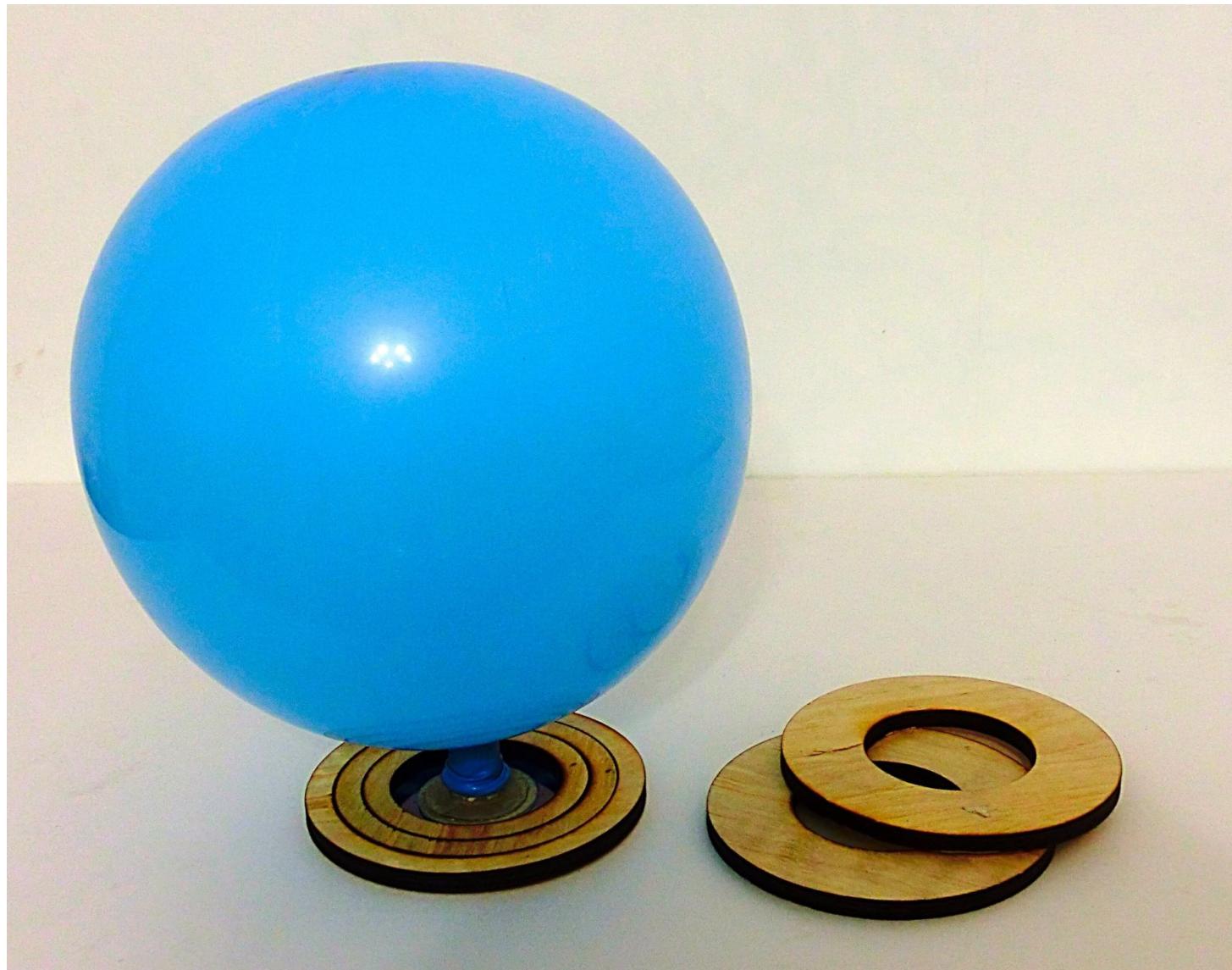
$$\tau = \frac{\tau_0}{\sqrt{1 - \frac{mg}{p \cdot \pi R^2} \cdot 2 \ln \frac{R}{a}}}$$

$$\tau = V \cdot \sqrt{\frac{9}{70} \cdot \frac{\rho \ln(R/a)}{2\pi a \cdot \pi R^2 \cdot \eta v_0}}$$

The hovering time almost does not depend on the weight of the vessel until this weight is not very large.

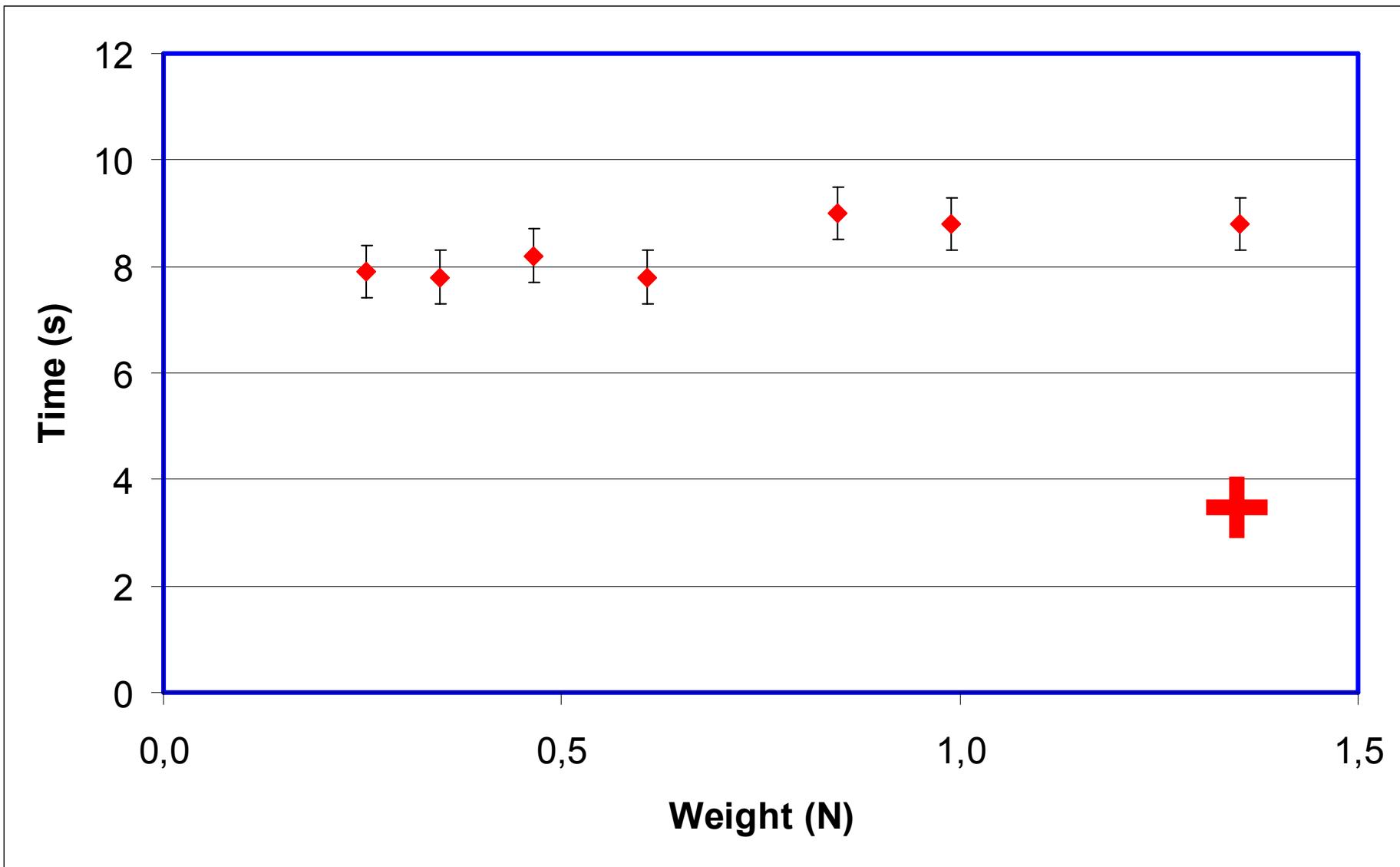
Experiment

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Time vs. weight (nozzle 13.8 mm)

42



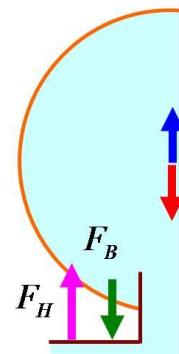
Summary

Conclusions

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Final force balance

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Pressure distribution under the disk

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Relative pressure (Pa)

Time (s)

Time vs. nozzle cross-section

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Hovering time with narrow nozzle

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For our

Hovering time with wide nozzle

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$$Q \cdot \delta \approx \frac{70}{9}$$

Theory-experiment comparison

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$\tau =$

Time vs. weight (nozzle 13.8 mm)

42

Time (s)

Weight (N)

- Jackson J.D., Symmons G.R. (1965) “An investigation of a laminar flow between two parallel disks”. *Appl. Sci. Res.* **15**, 59–75.
- Armengoll J., Calbó J., Pujol T., Roura P. (2011) “Bernoulli correction to viscous losses: Radial flow between two parallel discs”. *Am. J. Phys.* **76**, 730–737.
- Izarra Ch., Izarra G. (2014) “Stokes equation in a toy CD hovercraft”. *Eur. J. Phys.* **32**, 89–99.



**Thank you for
your attention!**